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# Fall 2024: Final Project

## COM-309: Quantum Information Processing

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This homework should be done by teams of two students. The implementation should be done with PennyLane **AND** Qiskit. Only one of the team's students should upload on Moodle a PDF with the answers to the theory questions and the two notebooks. The name of the two students and their SCIPER should be written in the PDF.

### 1 Werner state

The goal of this mini-project is to study the entanglement properties of mixed states. A mixed state  $\rho$  can be seen as a statistical mixture of pure states,

$$\rho = \sum_i p_i |\phi_i\rangle \langle \phi_i|, \quad p_i \in [0, 1], \quad \sum_i p_i = 1 \quad (1)$$

and can be represented by density matrices which are positive semi-definite, self-adjoint matrices of trace 1.

In this mini-project, we will focus on a specific class of states, called the Werner states. A Werner state  $\rho_W$  is a convex combination of the completely mixed state  $\mathbf{I}$  and the maximally entangled Bell state  $\rho_{11} = |B_{11}\rangle \langle B_{11}|$  where  $|B_{11}\rangle = \frac{|01\rangle - |10\rangle}{\sqrt{2}}$ . This mixture will depend on a parameter  $w \in [0, 1]$  such that

$$\rho_W(w) = (1 - w) \frac{\mathbf{I}}{4} + w \rho_{11}. \quad (2)$$

**Question Theory 1:** Give the matrix representation of  $\rho_W(w)$ , compute its eigenvalues, and check that  $\rho_W(w)$  is a valid state.

Werner state can be constructed with the following circuit

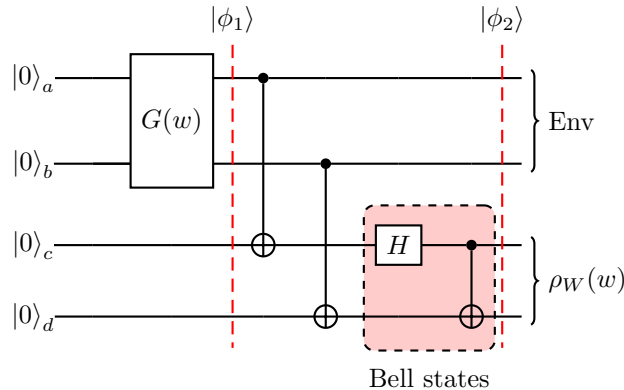


Figure 1: How to prepare Werner states

where  $G(w)$  is the unitary such that  $G(w) |00\rangle = \sqrt{\frac{1-w}{4}} (|00\rangle + |01\rangle + |10\rangle) + \sqrt{\frac{1+3w}{4}} |11\rangle$ . The two qubits  $a$  and  $b$  can be considered as the environment acting on our system, the two qubits  $c$  and  $d$ . In the red box in Figure 1, you can recognize the Bell state preparation.

**Question Theory 2:** Compute the state  $|\phi_2\rangle$  and check that  $\rho_W(w) = \text{Tr}_{a,b}(|\phi_2\rangle \langle \phi_2|)$ .

**Question Implementation 1:** Implement the circuit to construct Werner states in PennyLane and Qiskit. How to implement  $G(w)$  is given in the notebooks.

## 2 Separability and the Peres criterion

A pure state  $|\psi\rangle_{c,d}$  is a product state if it can be decomposed as the tensor product of  $|\psi_1\rangle_c$  and  $|\psi_2\rangle_d$  such that  $|\psi\rangle = |\psi_1\rangle_c \otimes |\psi_2\rangle_d$ . The state is otherwise entangled. This notion can be generalized to mixed states.

A mixed state  $\rho$  is *separable* if it can be written in the form

$$\rho_{c,d} = \sum_i p_i \rho_c^{(i)} \otimes \rho_d^{(i)}, \quad p_i \in [0,1], \quad \sum_i p_i = 1. \quad (3)$$

You can easily check that a product state is separable.

One can identify if a state is separable or not thanks to the **Peres criterion**: A two-qubit state  $\rho_{c,d}$  is entangled if and only if  $(\mathbf{I}_c \otimes T_d)\rho_{c,d}$  has a negative eigenvalue. Here,  $T_d$  is the operator that applies a transposition to the system  $d$ . In higher dimension, this criterion does not hold.

**Question Theory 3:** Prove that for a separable state  $\rho_{c,d}$ ,  $(\mathbf{I}_c \otimes T_d)\rho_{c,d}$  has only nonnegative eigenvalues. We will admit the other direction of the proof. *Hint 1:* What can you say about the eigenvalues of the transpose matrix? *Hint 2:* Remind that for a semi-positive definite matrix  $M$ , for any  $|\phi\rangle$ , we have  $\langle \phi | M | \phi \rangle \geq 0$ .

**Question Implementation 2:** Implement the Peres criterion to the Werner state for  $w \in [0,1]$ . For which  $w$  is  $\rho_W(w)$  separable?

**Question Theory 4:** Check your implementation by studying the eigenvalues of  $(\mathbf{I}_c \otimes T_d)\rho_W(w)$ .

## 3 Bell (CHSH) inequality

In the lectures, we discussed the Bell inequalities. They are derived in relation to the experiments on quantum systems: if the results of an experiment can be explained by the local hidden-variable (LHV) model, such inequality should be satisfied. Its violation shows that in fact, LHV theory does not hold. Particularly, we looked at the CHSH inequality: in the corresponding experiment, Alice and Bob shared the two-qubit (pure) *Bell* state  $|\Psi\rangle$  and estimated the correlation coefficient  $\chi = \langle \Psi | \mathcal{B} | \Psi \rangle$  by measuring the observable

$$\mathcal{B} = A \otimes B + A \otimes B' - A' \otimes B + A' \otimes B',$$

where  $A, A', B, B'$  are observables

$$\begin{aligned} A &= |\alpha\rangle \langle \alpha| - |\alpha_\perp\rangle \langle \alpha_\perp|, & A' &= |\alpha'\rangle \langle \alpha'| - |\alpha'_\perp\rangle \langle \alpha'_\perp|, \\ B &= |\beta\rangle \langle \beta| - |\beta_\perp\rangle \langle \beta_\perp|, & B' &= |\beta'\rangle \langle \beta'| - |\beta'_\perp\rangle \langle \beta'_\perp|. \end{aligned}$$

and  $\{|\alpha\rangle, |\alpha_\perp\rangle\}$ ,  $\{|\alpha'\rangle, |\alpha'_\perp\rangle\}$ ,  $\{|\beta\rangle, |\beta_\perp\rangle\}$ ,  $\{|\beta'\rangle, |\beta'_\perp\rangle\}$  are pairs of orthogonal states. The CHSH inequality states that  $|\chi_{\text{LHV}}| \leq 2$ , but in practice, with certain choice of configuration  $\alpha, \alpha', \beta, \beta'$ ,  $|\chi_{\text{QM}}|$  can reach up to  $2\sqrt{2}$ ! In fact, one can prove that for any two-qubit pure *entangled* state  $|\Psi\rangle$  there exists such configuration that the CHSH inequality is violated.

However, now we consider mixed states as well, so the correlation coefficient is now expressed as

$$\chi = \text{Tr}(\rho \mathcal{B}).$$

**Question Theory 5:** In homework 6, you have shown that the inequality  $\langle \Psi | \mathcal{B} | \Psi \rangle \leq 2\sqrt{2}$  holds for *any pure state*  $|\Psi\rangle \in \mathbb{C}^2 \otimes \mathbb{C}^2$  (Tsirelson's bound). i.e.  $\chi(w) \leq 2\sqrt{2}$ , holds. Using this, show that it holds for arbitrary two-qubit mixed state as well, i.e.  $\chi(w) \leq 2\sqrt{2}$ .

One can show that if the mixed state  $\rho$  is separable, the CHSH inequality is satisfied. This does not work in reverse: mixed non-separable states can satisfy CHSH inequality for any configuration. In this problem, we will see this with the example of Werner states, i.e. we will measure correlation coefficient

$$\chi(w) = \text{Tr}(\rho_W(w) \mathcal{B}).$$

**Question Theory 6:** Show that the optimal configuration  $\alpha, \alpha', \beta, \beta'$  does not depend on value of  $w$ .

One can prove that  $\alpha = \frac{\pi}{4}, \alpha' = 0, \beta = -\frac{\pi}{8}, \beta' = -\frac{3\pi}{8}$  is an optimal configuration.

**Question Theory 7:** For this configuration, express  $A, A', B, B'$  as a linear combination of Pauli matrices. This will be used later for PennyLane and Qiskit implementations of observables.

**Question Implementation 3:** Run the circuits and measure  $\mathcal{B}$  to evaluate  $\chi(w)$  for  $w \in [0, 1]$ . Draw a plot of this dependence; compare it to 2 (validity of the CHSH inequality bound) and Tsirelson's bound. More details are provided in the corresponding notebooks.

**Question Theory 8:** From the implementation, what is the value of  $\chi(0)$  and  $\chi(1)$ ? What are the values of  $w$  s.t. the CHSH inequality is satisfied? Compare them with the values of  $w$  for those when  $\rho_W(w)$  is separable.