

The Schmidt Theorem.

①

Theorem Let $|\psi\rangle \in \mathcal{H}_A \otimes \mathcal{H}_B$ be a pure state of a bipartite system. Let $\rho_{AB} = |\psi\rangle\langle\psi|$ the corresponding density matrix.

Then $\rho_A = \text{Tr}_B \rho_{AB}$ and $\rho_B = \text{Tr}_A \rho_{AB}$ have the same non-zero eigenvalues.

Moreover $S(\rho_A) = S(\rho_B)$ since of $\{\lambda_\alpha^A\}$

& $\{\lambda_\beta^B\}$ are the set of non zero eigenvalues of ρ_A & ρ_B

we have $S(\rho_A) = - \sum_\alpha \lambda_\alpha^A \ln \lambda_\alpha^A$; $S(\rho_B) = - \sum_\beta \lambda_\beta^B \ln \lambda_\beta^B$.

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Proof. Based on SVD decomposition.

$$\text{let } |\psi\rangle = \sum_{i=1}^{d_A} \sum_{j=1}^{d_B} \psi_{ij} |i\rangle_A \otimes |j\rangle_B$$

where $d_A = \dim \mathcal{H}_A$, $d_B = \dim \mathcal{H}_B$.

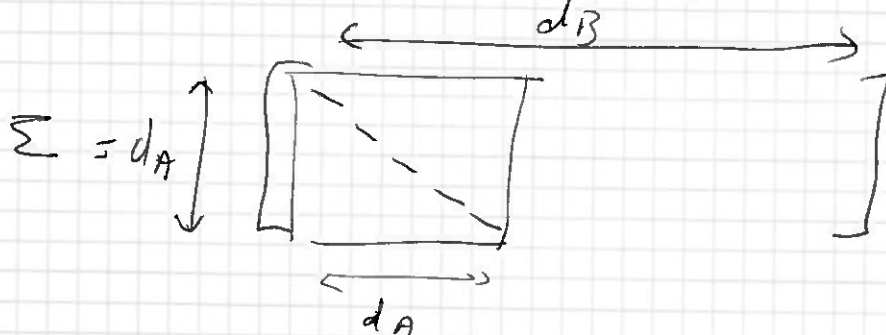
Assume the case $d_A \leq d_B$ (other one is similar).

By SVD the "matrix" ψ_{ij} which is $d_A \times d_B$:

$$\psi_{ij} = \sum_{k=1}^{d_A} \sum_{l=1}^{d_B} U_{ik} \Sigma_{kl} V_{jl}^*$$

with $U : d_A \times d_A$ unitary, $V : d_B \times d_B$ unitary

$\Sigma : d_A \times d_B$ a ^(diagonal) matrix of sing values $\sigma_\alpha \geq 0$.



We have

$$|\psi\rangle \langle \psi| = \sum_{k=1}^{d_A} \sum_{l=1}^{d_B} \sum_{k'=1}^{d_A} \sum_{l'=1}^{d_B} \psi_{ij} \psi_{k'l}^* |i, j\rangle_{AB} \langle k', l|_{AB}$$

The density matrix $\rho = |\psi\rangle\langle\psi|$ has matrix elements

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$$\rho_{ij;lm} = \psi_{ij} \psi_{lm}^*$$

Thus $\rho_A = \text{Tr}_B \rho_{AB}$ has matrix elements

$$(\rho_A)_{ie} = \sum_{j=1}^{d_B} \psi_{ij} \psi_{ej}^*$$

$$\Rightarrow \boxed{\rho_A = \psi \psi^\dagger} \quad \text{where } \psi \text{ is matrix } \psi_{ij}$$

Now the SVD is $\psi = U \Sigma V^\dagger$

$$\Rightarrow \rho_A = U \Sigma V^\dagger V \Sigma^\dagger U^\dagger$$

$$= U \underbrace{\Sigma \Sigma^\dagger}_{d_A \times d_A \text{ matrix}} U^\dagger$$

since $V^\dagger V = \mathbb{1}_{d_B \times d_B}$

$d_A \times d_A$ matrix

$$\Sigma \Sigma^\dagger = \left[\begin{array}{c} \boxed{\text{diagonal}} \\ \vdots \end{array} \right] \left[\begin{array}{c} \boxed{\text{diagonal}} \\ \vdots \end{array} \right]$$

$$= \boxed{\text{diagonal}}^2 \quad \{\text{matrix of } \sigma_i^2\}$$

On the other hand:

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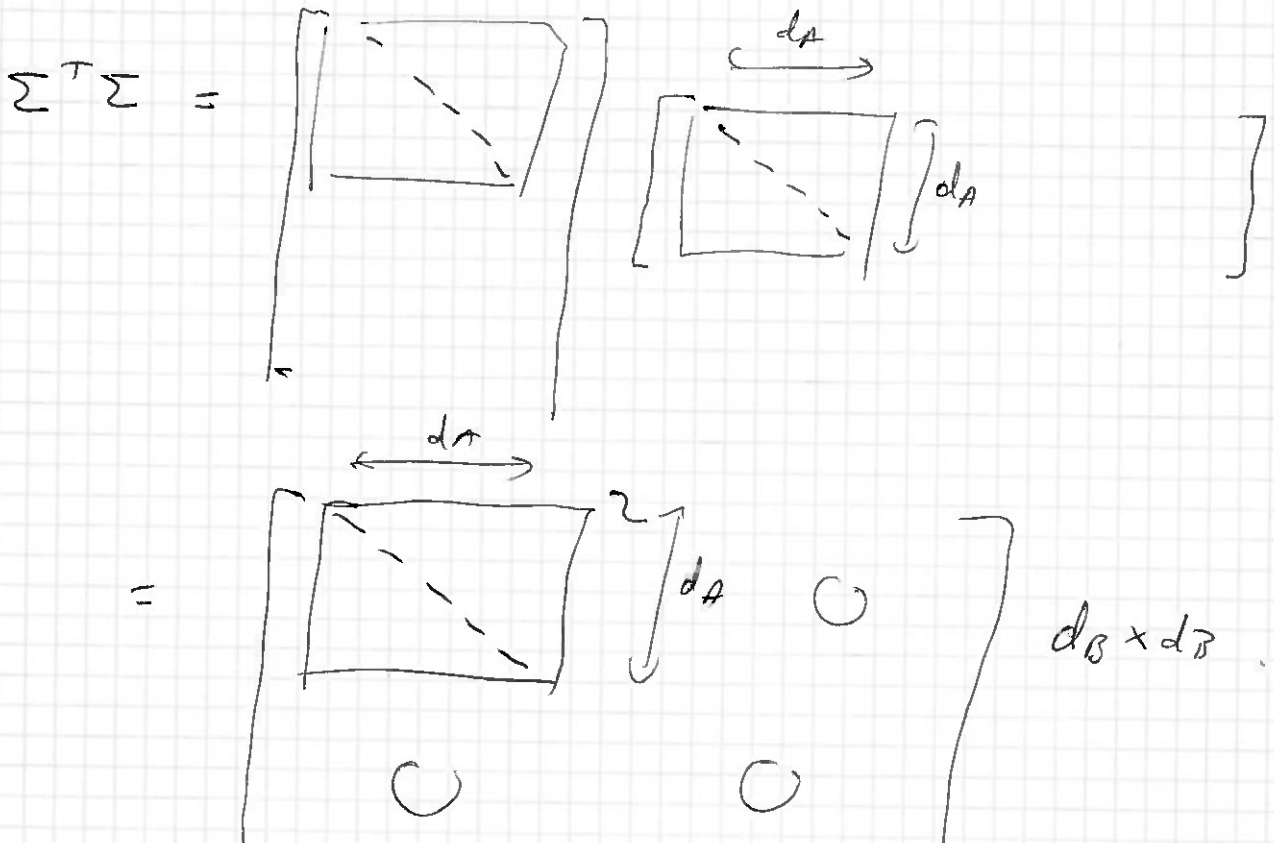
$\rho_B = \text{Tr}_A \rho_{AB}$ has matrix elements:

$$(\rho_B)_{jm} = \sum_{i=1}^{d_A} \psi_{ij} \psi_{im}^* = \sum_{i=1}^{d_A} (\psi^\dagger)_{mi} \psi_{ij}$$

$$\Rightarrow \boxed{\rho_B = \psi^\dagger \psi}$$

Now from $\psi = U \Sigma V^\dagger$ we get

$$\rho_B = V \underbrace{\Sigma^\dagger U^\dagger U}_{\substack{\uparrow \\ d_A \times d_A}} \Sigma V^\dagger = V \underbrace{\Sigma^\dagger \Sigma}_{\substack{\uparrow \\ d_B \times d_B}} V^\dagger$$



has non zero eigenvalues given by $\left[\begin{array}{c} \diagup \\ \text{---} \\ \diagdown \end{array} \right]^2$ same as ρ_A (CFD)