

# Quantum computation: lecture ~~10~~ 9

## Grover's algorithm

Let  $f: \{0,1\}^n \rightarrow \{0,1\}$  be a Boolean function  
and  $A = \{x \in \{0,1\}^n : f(x) = 1\}$

We are considering the search problem,  
namely to identify one element  $x$  of  $A$   
with as few calls as possible to the oracle  $f$ .

Consider first the case where  $A = \{x^*\}$ , singleton:

Classically, identifying  $x^*$  may require up to  $N = 2^n$  calls to the oracle  $f$ , in the worst case.

As we will see, Grover's quantum algorithm only requires  $\sqrt{N} = 2^{n/2}$  calls to the oracle  $f$  to identify  $x^*$ .

Note: Grover's algorithm is usually presented solving the following problem: given a directory of  $N$  names in alphabetical order and corresponding phone numbers, it allows to recover the name corresponding to a given phone number in only  $\sqrt{N}$  steps (instead of  $N$  steps classically).

However, in order to work, Grover's algorithm requires us to build the gate  $U_f$ , which is not doable in the directory search pb, as we know nothing about the function  $f$ !

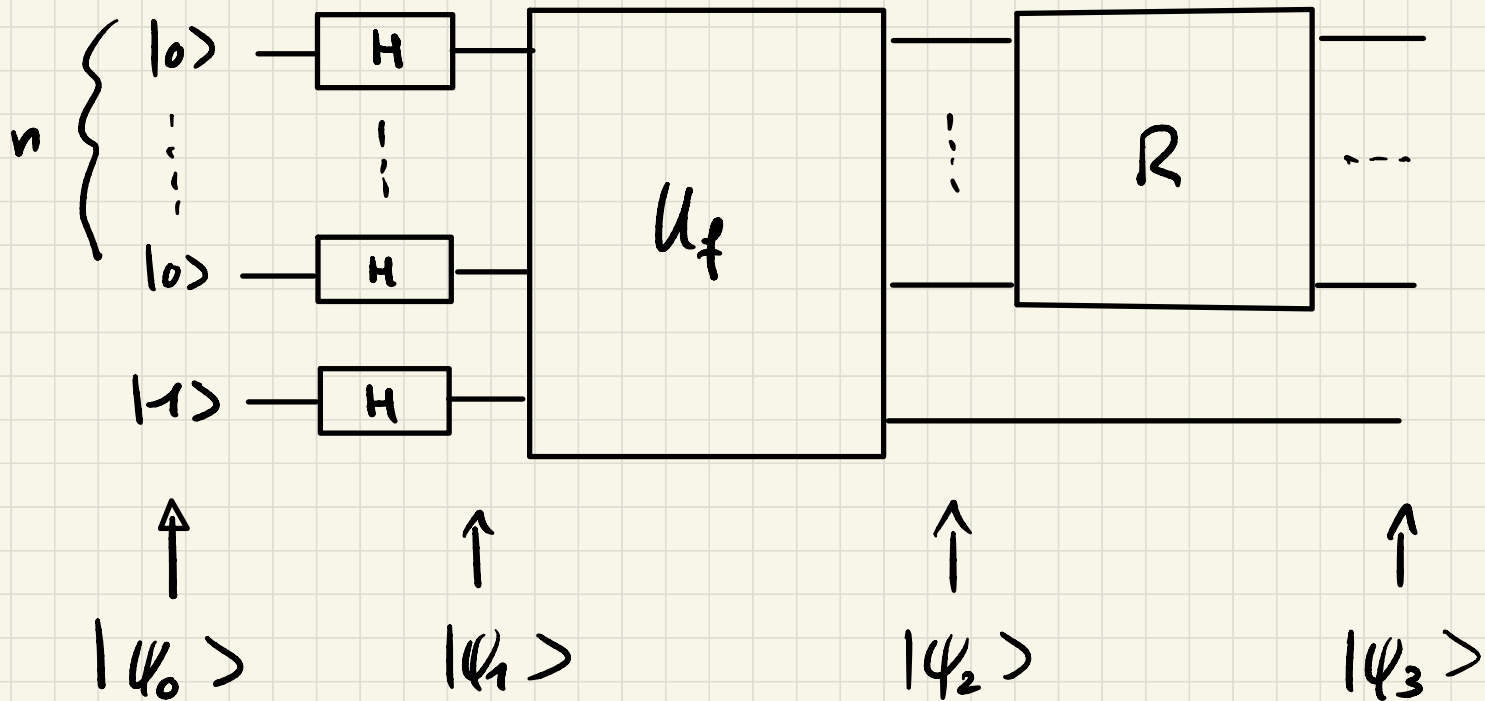
But please be patient: we will see later other interesting applications of Grover's algorithm.

Note also that we will consider in general functions  $f$  with  $|A| = M \in \{1..N\}$ .

Surprisingly perhaps, considering this generalization (without focusing on the case  $M=1$ ) will help us visualize better how the algorithm works!

# Grover's quantum circuit

"reflection gate"  
↓



Let us compute (as already done multiple times)

$$\bullet \quad |\psi_1\rangle = H|0\rangle \otimes \dots \otimes H|0\rangle \otimes H|1\rangle$$

$$= \frac{1}{\sqrt{N}} \sum_{x \in \{0,1\}^n} |x\rangle \otimes \underline{|-\rangle}$$

$$\left[ = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle) \right]$$

$$\bullet \quad |\psi_2\rangle = \frac{1}{\sqrt{N}} \sum_{x \in \{0,1\}^n} (-1)^{f(x)} |x\rangle \otimes |-\rangle$$

So far, nothing new. Now, remember that in our case,  $\Pi = |A| = \{x : f(x) = 1\}$ , so  $N - \Pi = |A^c| = \{x : f(x) = 0\}$ .  
Therefore:

$$|\psi_2\rangle = \frac{1}{\sqrt{N}} \sum_{x \in \{0,1\}^n} (-1)^{f(x)} |x\rangle \otimes |-\rangle$$

$$= \frac{1}{\sqrt{N}} \left( \sqrt{\frac{N-\Pi}{N-\Pi}} \sum_{x \in A^c} |x\rangle - \sqrt{\frac{\Pi}{\textcolor{red}{N}}} \sum_{x \in A} |x\rangle \right) \otimes |-\rangle$$

$$= \left( \sqrt{\frac{N-\Pi}{N}} \left( \frac{1}{\sqrt{N-\Pi}} \sum_{x \in A^c} |x\rangle \right) - \sqrt{\frac{\Pi}{N}} \left( \frac{1}{\sqrt{\Pi}} \sum_{x \in A} |x\rangle \right) \right) \otimes |-\rangle$$

Let us write  $|P\rangle = \frac{1}{\sqrt{N-M}} \sum_{x \in A^c} |x\rangle$

and  $|S\rangle = \frac{1}{\sqrt{M}} \sum_{x \in A} |x\rangle$  : both  $|P\rangle$  and  $|S\rangle$

are quantum states (normalized to 1)

and  $|\psi_2\rangle$  can be rewritten as

$$|\psi_2\rangle = \left( \sqrt{\frac{N-M}{N}} |P\rangle - \sqrt{\frac{M}{N}} |S\rangle \right) \otimes \underset{\uparrow}{|-\rangle}$$

[From now on, we will forget the extra  $|-\rangle$  state.]

Note also that  $\left(\sqrt{\frac{N-M}{N}}\right)^2 + \left(\sqrt{\frac{M}{N}}\right)^2 = \frac{N-M}{N} + \frac{M}{N} = 1$ ,

so there exists  $\Theta_0 \in [0, \frac{\pi}{2}]$  such that

$$\cos \Theta_0 = \sqrt{\frac{N-M}{N}} \quad \text{and} \quad \sin \Theta_0 = \sqrt{\frac{M}{N}}$$

$$\text{and } |\psi_2\rangle = \cos \Theta_0 |P\rangle - \sin \Theta_0 |S\rangle$$

Likewise, observe that

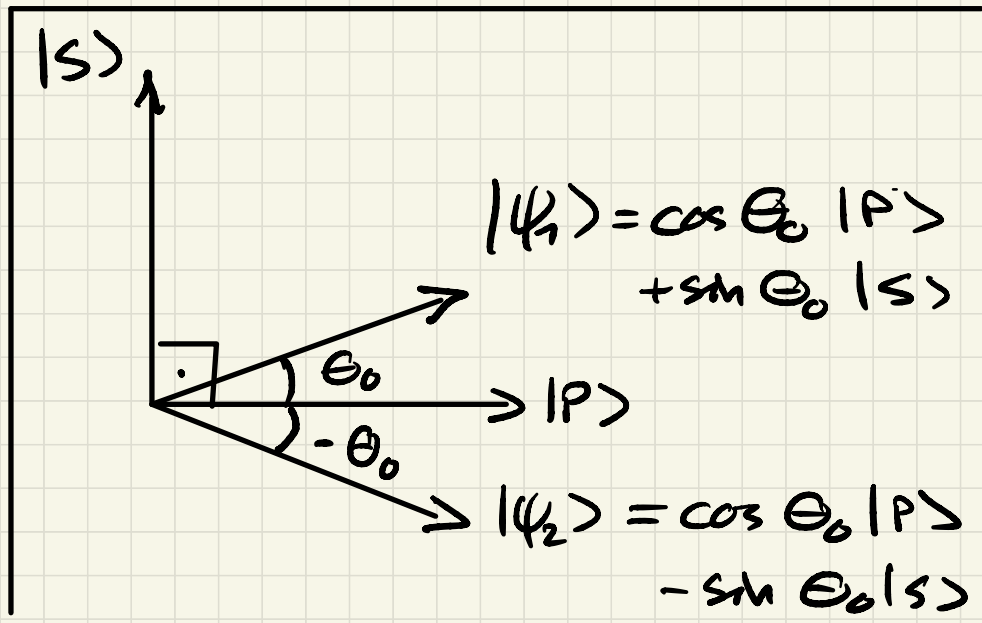
$$|\psi_1\rangle = \cos \Theta_0 |P\rangle + \sin \Theta_0 |S\rangle$$

(if we again forget the extra  $|- \rangle$  state)

# Geometric interpretation

$$|P\rangle = \frac{1}{\sqrt{N-M}} \sum_{x \in A^c} |x\rangle \quad \text{and} \quad |S\rangle = \frac{1}{\sqrt{M}} \sum_{x \in A} |x\rangle$$

are orthogonal (as they share no common basis element), so we obtain the following picture:

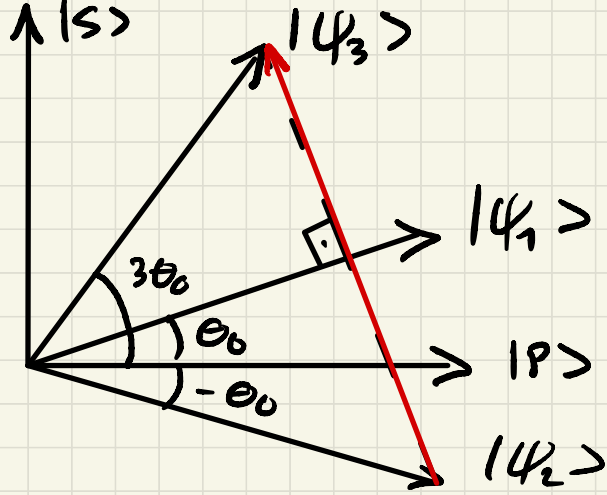


The action of the gate  $U_f$  on state  $|\psi_1\rangle$  can therefore be interpreted as a reflection with respect to the axis  $|P\rangle$  !

But note that we do not know the axes  $|P\rangle$  and  $|S\rangle$ ; this is exactly what we are after; more precisely, our aim now is to push as much as possible the state of the system towards  $|S\rangle$ , which contains only elements  $x \in A$ .

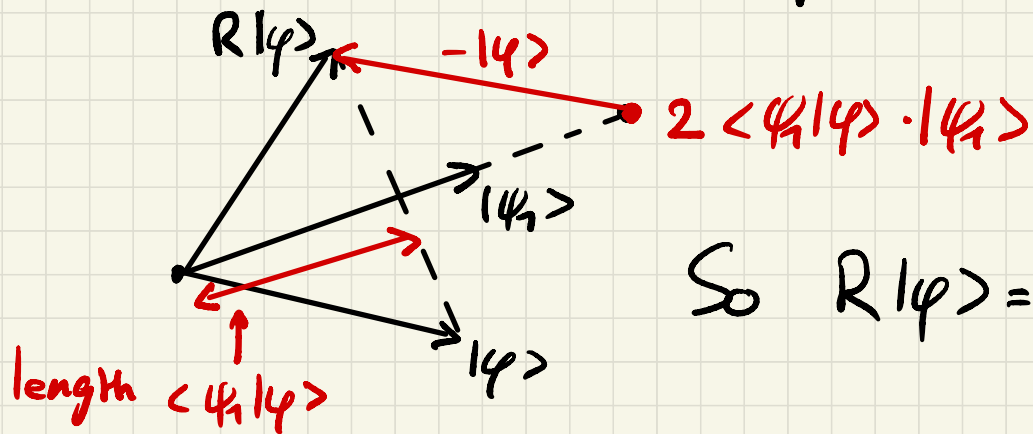
## Reflection gate R

A first step in this direction is done by applying the gate R, which is another reflection with respect to state  $|\psi_1\rangle$ :



Building such a gate  $R$  does not require using the function  $f$  again, as remember that  $|\psi_1\rangle = \frac{1}{\sqrt{N}} \sum_{x \in \{0,1\}^n} |x\rangle$  (~~0~~  $1 \rightarrow$ ), simply.

Here is the geometric procedure to build  $R$ :



$$\text{So } R|\psi\rangle = 2 \langle \psi_1 | \psi \rangle \cdot |\psi_1\rangle - |\psi\rangle$$

$$\text{We have } R|\varphi\rangle = 2\langle\varphi_1|\varphi\rangle \cdot |\varphi_1\rangle - |\varphi\rangle$$

$$= 2|\varphi_1\rangle\langle\varphi_1|\varphi\rangle - |\varphi\rangle$$

$$= \underbrace{(2|\varphi_1\rangle\langle\varphi_1| - I_n)}_{\text{matrix!}} \cdot |\varphi\rangle$$

$$= (2H^{\otimes n}|\varphi_0\rangle\langle\varphi_0|H^{\otimes n} - I_n) \cdot |\varphi\rangle$$

$$= H^{\otimes n} \underbrace{(2|\varphi_0\rangle\langle\varphi_0| - I_n)}_{\text{matrix!}} H^{\otimes n} \cdot |\varphi\rangle$$

$$= |0\rangle \otimes \dots \otimes |0\rangle \quad (\text{we forget again state } |-\rangle)$$

## Exercise:

Build the gate  $2|\psi_0\rangle\langle\psi_0| - I_n$

Hint: observe that

$$(2|\psi_0\rangle\langle\psi_0| - I_n)|\varphi\rangle$$

$$= |\varphi\rangle \quad \text{if } |\varphi\rangle = |\psi_0\rangle = |0\rangle \otimes \dots \otimes |0\rangle$$

$$\begin{cases} -|\varphi\rangle & \text{if } |\varphi\rangle \text{ is any other basis element} \end{cases}$$

Now, how to proceed from there on?

- With the successive application of  $U_f$  and  $R$ , the state evolves from

$$|\psi_1\rangle = \cos \theta_0 |P\rangle + \sin \theta_0 |S\rangle \quad \text{angle } +\theta_0$$

to

$$|\psi_2\rangle = \cos \theta_0 |P\rangle - \sin \theta_0 |S\rangle \quad \text{angle } -\theta_0$$

to

$$|\psi_3\rangle = \cos(3\theta_0) |P\rangle + \sin(3\theta_0) |S\rangle \quad \text{angle } +3\theta_0$$

Therefore, the successive application of  $U_f$  and  $R$  corresponds to a rotation of angle  $+2\theta_0$ , which brings the state closer to state  $|s\rangle$ , which is our aim. By iterating this operation an appropriate number of times, we can get arbitrarily close to  $|s\rangle$   $\rightarrow$  next week!