

Nuclear Fusion and Plasma Physics

Lecture 12

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Wrapping it up....
**Burning plasmas, ITER safety, and
route to a fusion power plant**

Burning plasma: generalities

Fast ion loss mechanisms

Burn stability and control

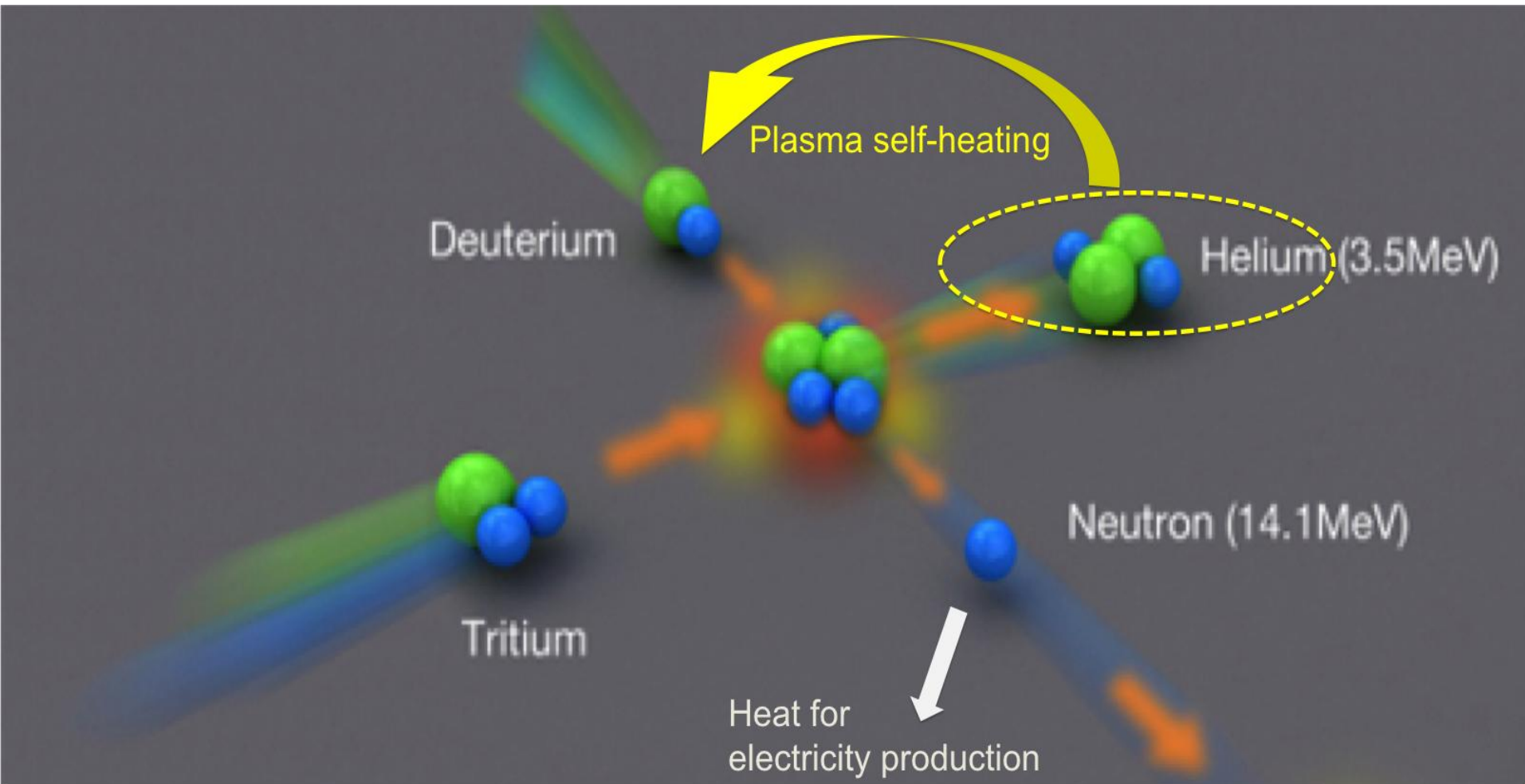
Tritium in fusion

ITER safety and licensing

Toward a fusion power plant

BURNING PLASMAS: GENERALITIES

D-T fusion reaction



Definition of a burning plasma

Fusion power density $\equiv P_{\text{fusion}} = \frac{1}{4}n^2 \langle \sigma v \rangle E_{\text{fusion}} \quad (n_D = n_T = n/2)$

α power density $\equiv P_{\alpha} = 0.2 P_{\text{fusion}}$

Thermal energy density $W \equiv 3nT$

Energy balance

$$dW/dt = P_{\alpha} + P_{\text{in}}/V - W/\tau_E$$

α -heating
 ext. heating
 losses

confinement
time



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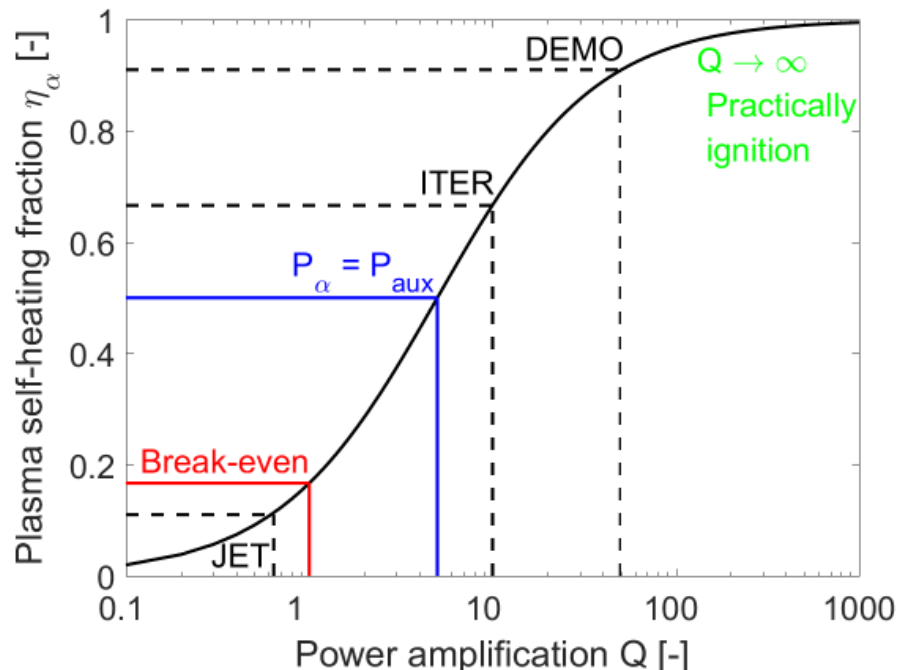
$$dW/dt = P_\alpha + P_{\text{in}}/V - W/\tau_E$$

α -heating *ext. heating* *losses*

confinement
time

$$Q \equiv P_{\text{fusion}}/P_{\text{in}} = 5 P_\alpha/P_{\text{in}}$$

$$\alpha \text{ heating fraction: } f_\alpha \equiv P_\alpha/(P_\alpha + P_{\text{in}}) = Q/(Q+5)$$

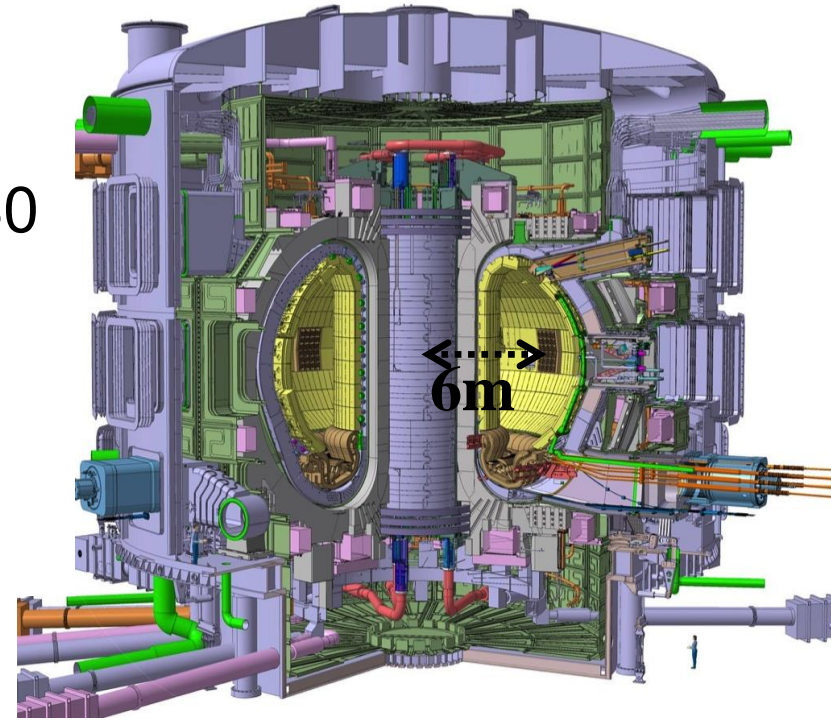


$Q \geq 10$; $P_{\text{fusion}} \geq 500\text{MW}$; $\sim 500\text{s}$

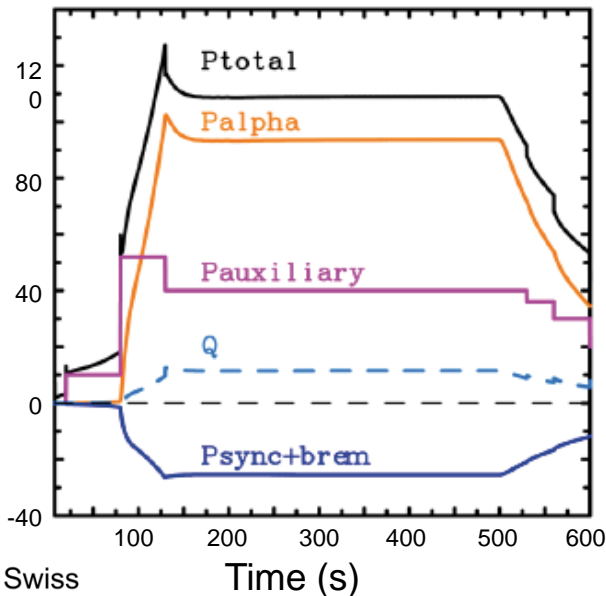
Explore steady-state $Q \geq 5$, $\leq 3000\text{s}$

May explore 'controlled ignition' $Q \geq 30$

→ burning plasma



Y.Gribov



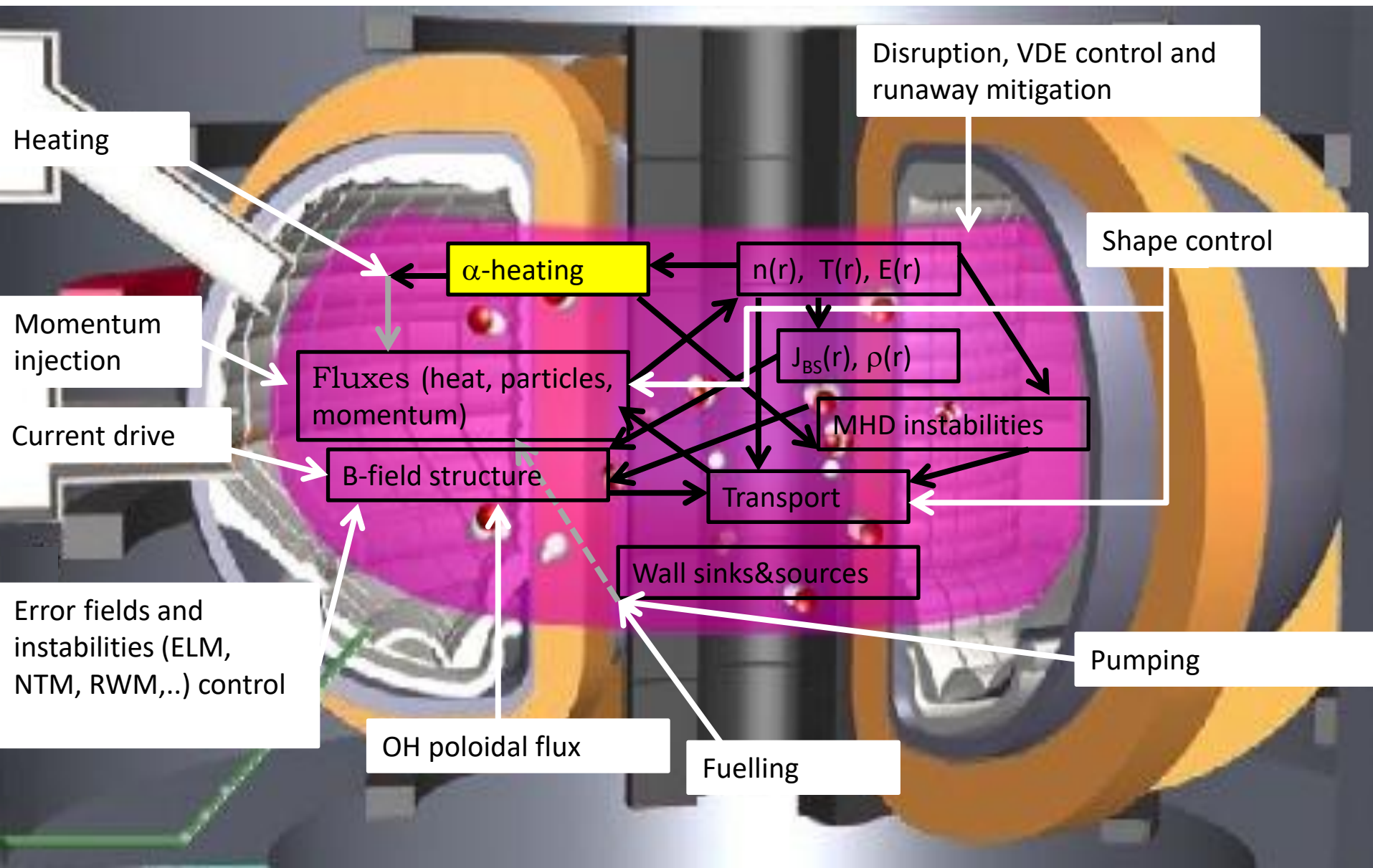
Q=10 scenario

$I_p = 15\text{ MA}$, $P_{\text{fus}} \sim 500\text{ MW}$

Total α energy in single shot: $\sim 50\text{GJ}$

Total α energy produced by JET and TFTR DT campaigns: $< 10\text{GJ}$

Burning plasmas: interplay of plasma dynamics and external systems



High performance

Operational limits, heat flux on plasma facing components

Nuclear environment

Radiation T retention, dust, breeding

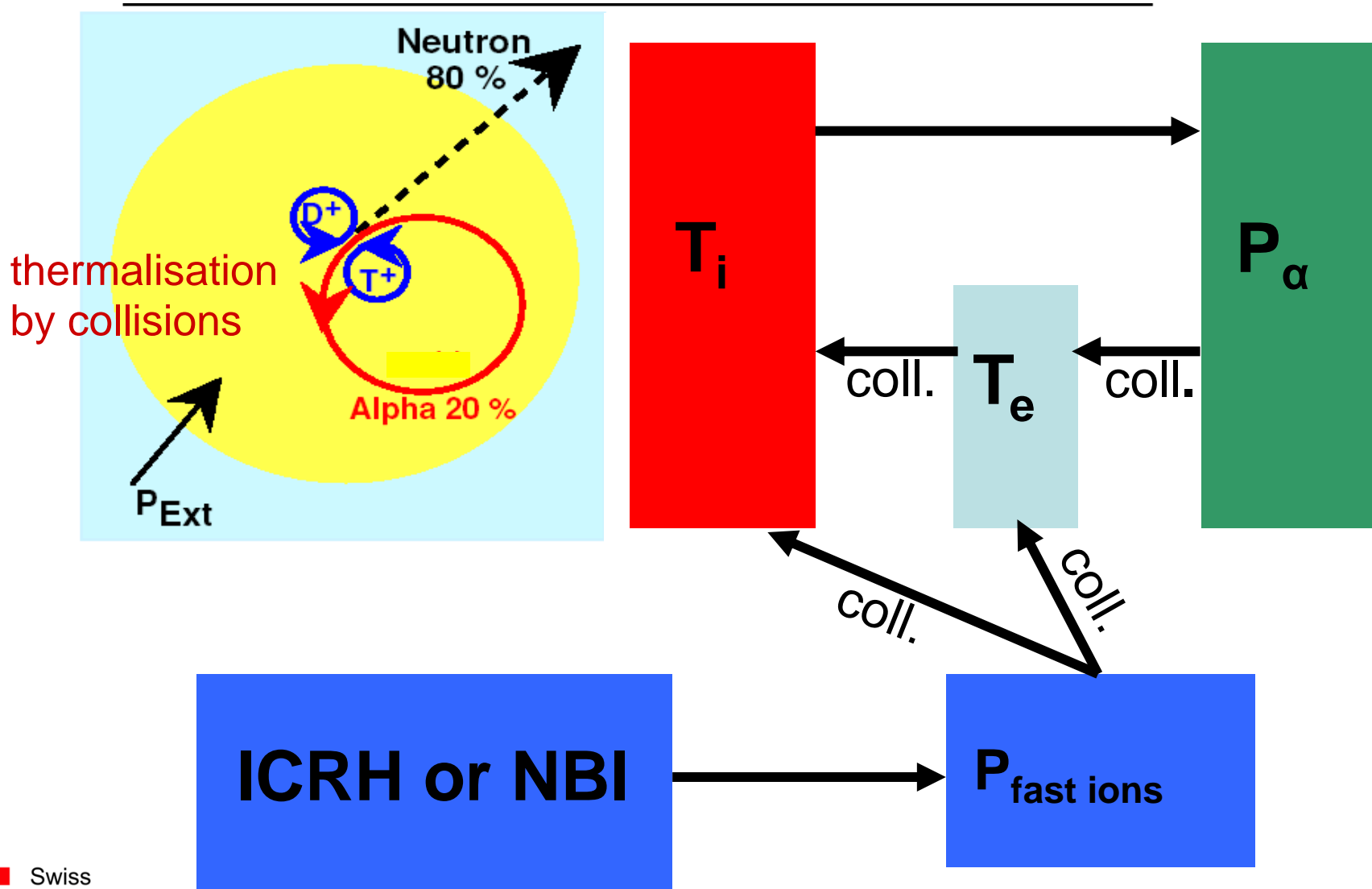
New physics parameter range

Small ρ_{Li}/a , high density and high temperature, low collisionality

Large isotropic population of fusion produced α 's, dominant source of plasma heating

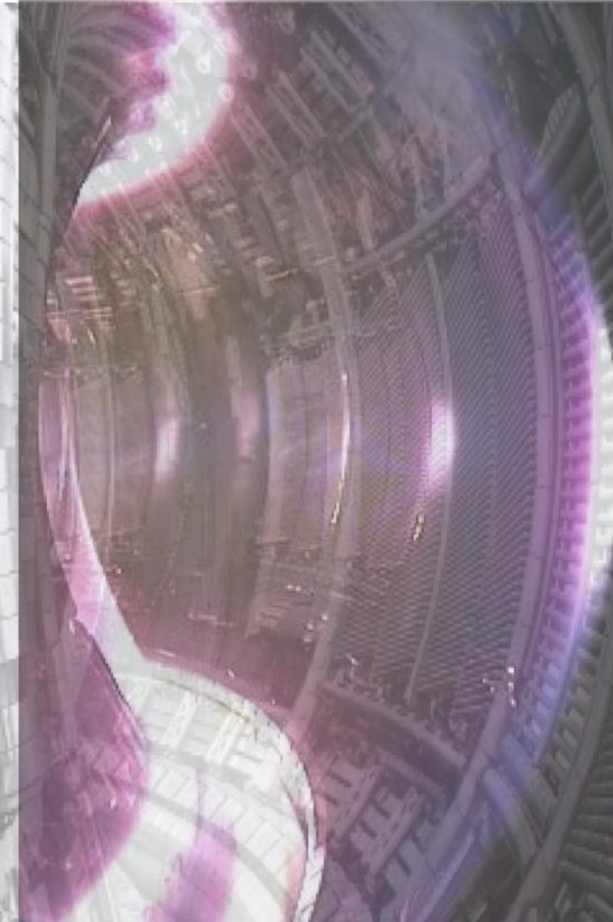
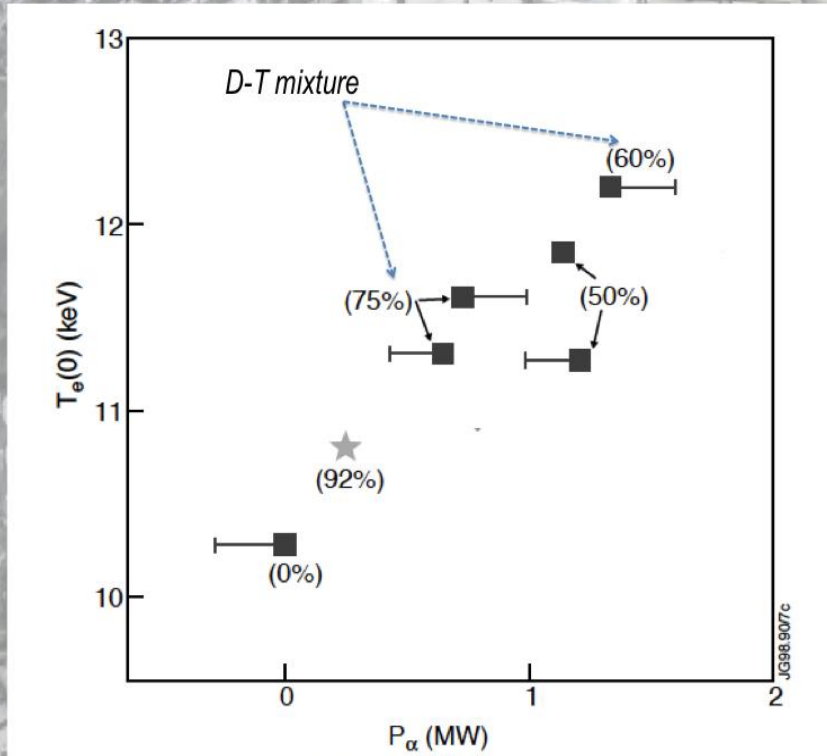
Coupling of profiles $P_{\alpha}(r)$, $p(r)$, $q(r)$, $E_r(r)$, diffusion coefficient(r), $n_{He}(r)$, ...

External heating and self-heating



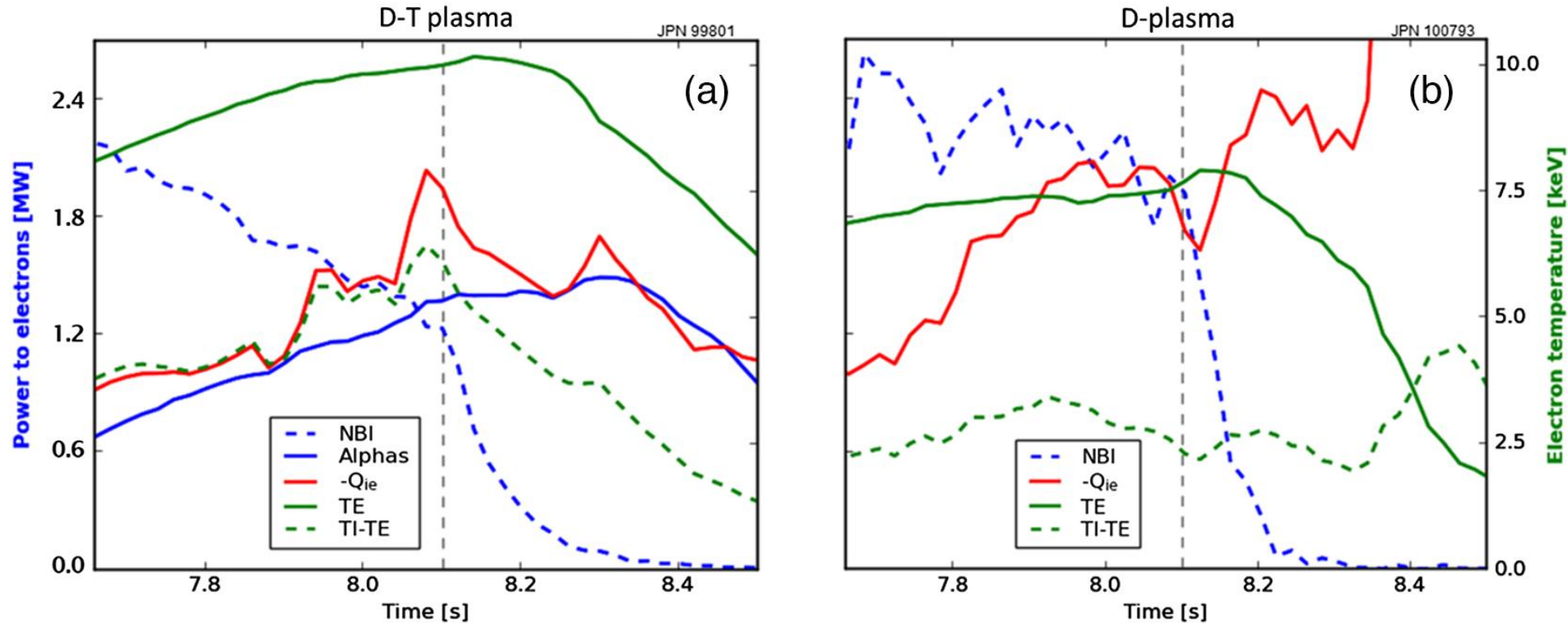
EPFL Fast ions and the self-heating process

Electron heating by fusion α 's - 1997



EPFL Fast ions and the self-heating process

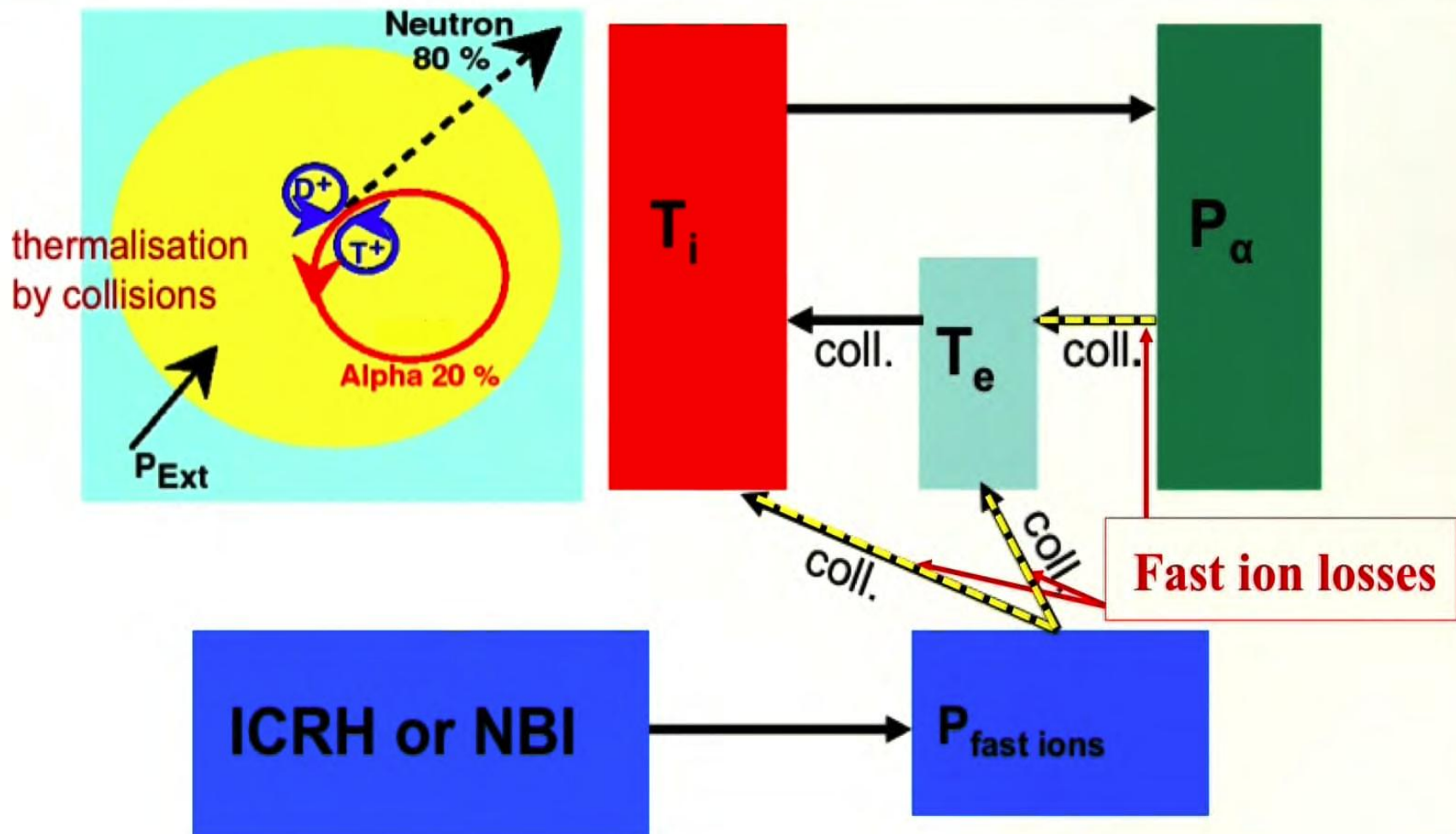
Electron heating by fusion α 's - 2023



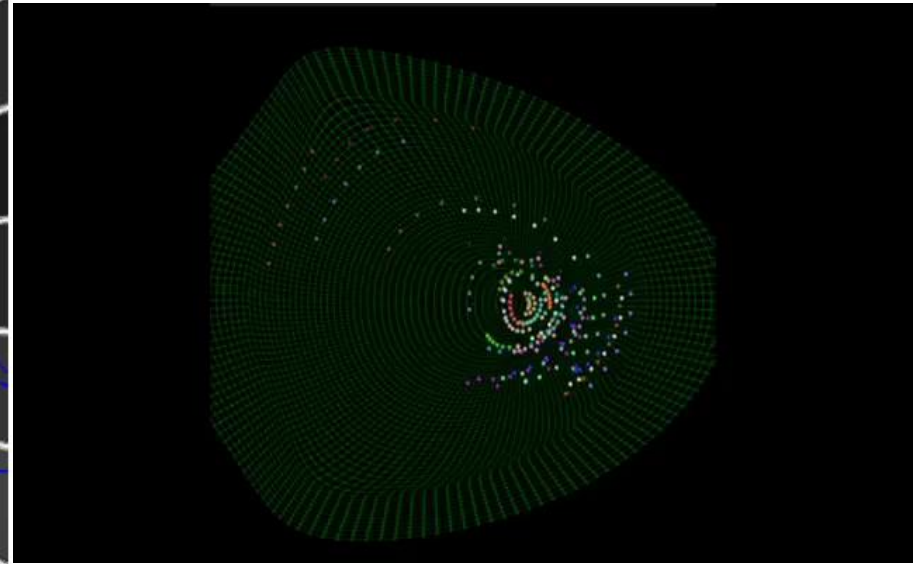
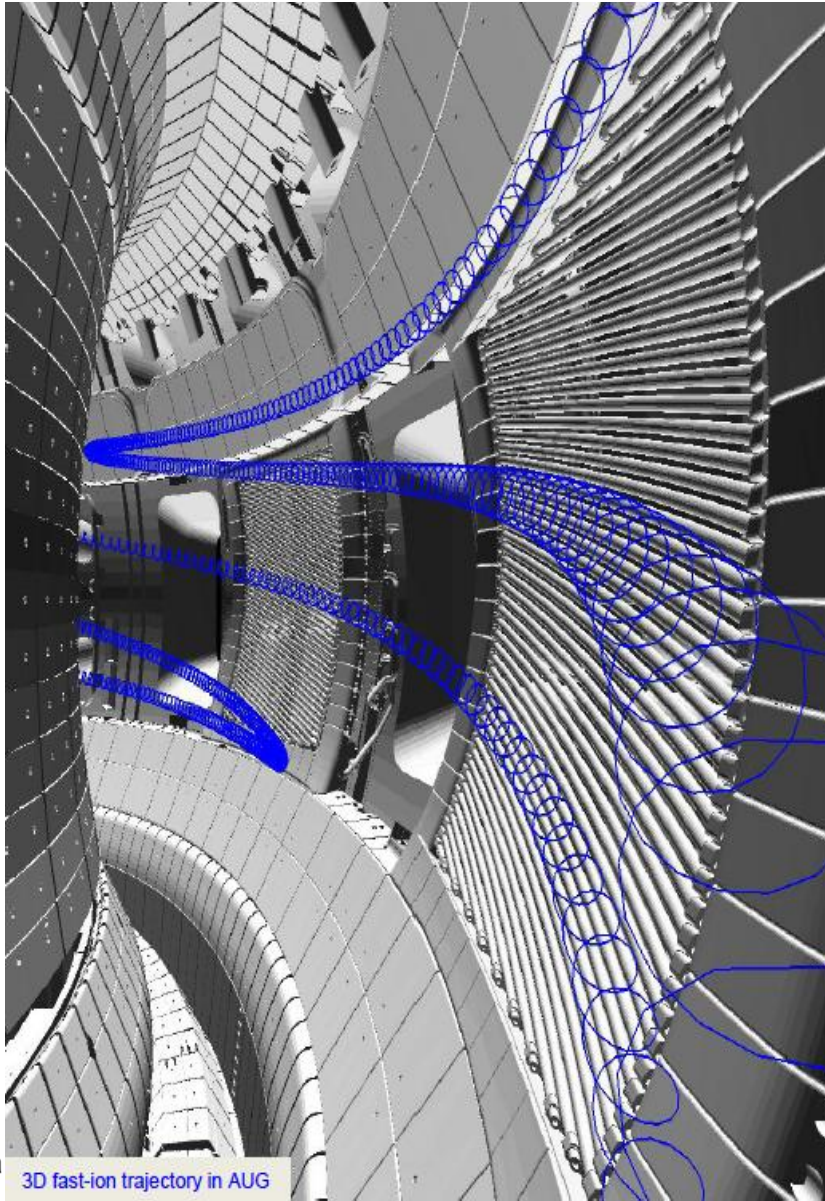
TRANSP [10] analysis of electron heating in JET D-T discharge No. 99801(a) and deuterium discharge No. 100793 (b). The power transferred to electrons (left scale) by alphas, NBI, and thermal ions ($-Q_{ie}$) are presented as well as electron temperature on axis (TE) and a difference between the ion and electron temperatures (TI-TE) in the plasma core (right scale); dashed line marks the NBI power cut.

External heating and self-heating

To reach and sustain burning plasma regime, fast ions must be well confined



To reach and sustain burning plasma regime, fast ions must be well confined



GPU-NUBEAM DIII-D
Simulation

To reach and sustain burning plasma regime, fast ions must be well confined

Need to understand and possibly minimize redistribution and loss mechanisms

- Magnetic field imperfections

- Low frequency MHD instabilities

- Turbulence

- Resonant interaction with Alfvén waves

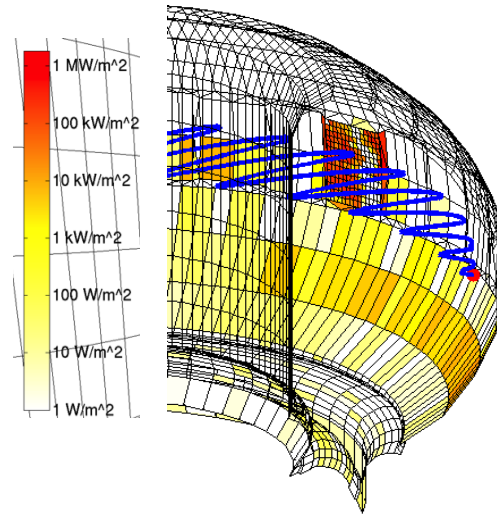
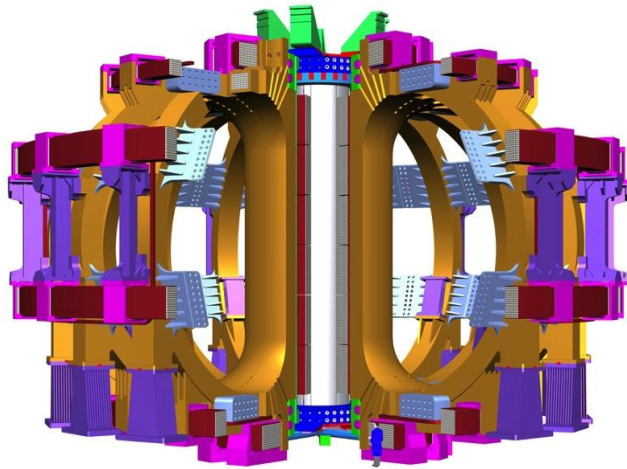
FAST ION LOSS MECHANISMS

MAGNETIC FIELD IMPERFECTIONS

EPFL Fast ions and the self-heating process

Fast ion losses in 3D B-field

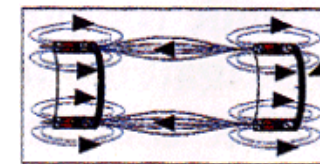
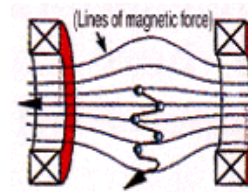
B-field inhomogeneities can lead to orbit trapping and losses



ASCOT code
T.Kurki-Suonio

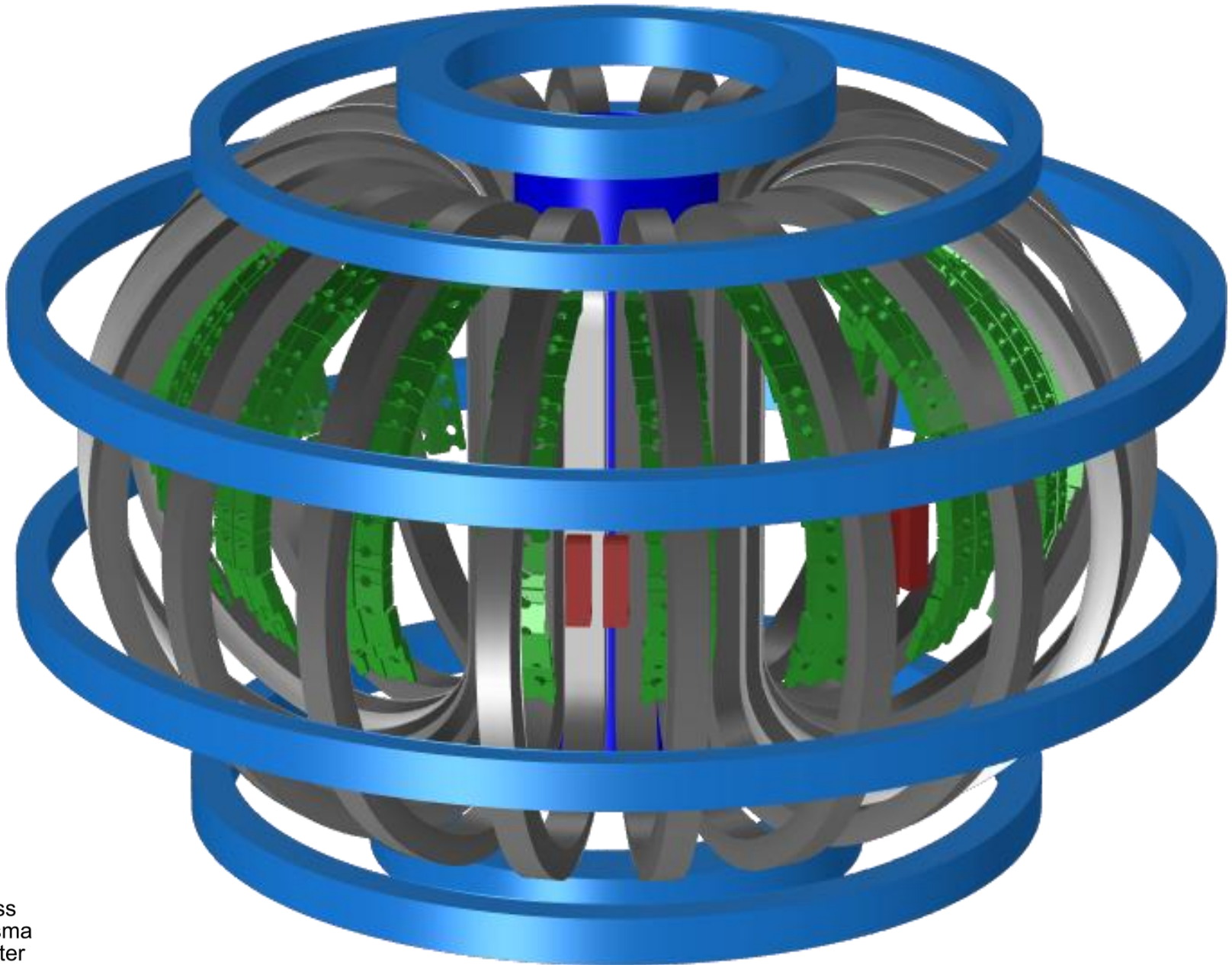
Ripple effects on fast ions qualitatively understood

Ferritic inserts in ITER to reduce ripple losses (from ~6% to ~0.4%)



Ferritic steel
(Ferromagnetic material)

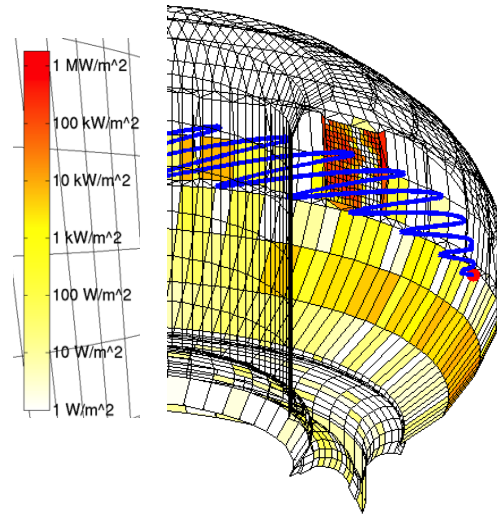
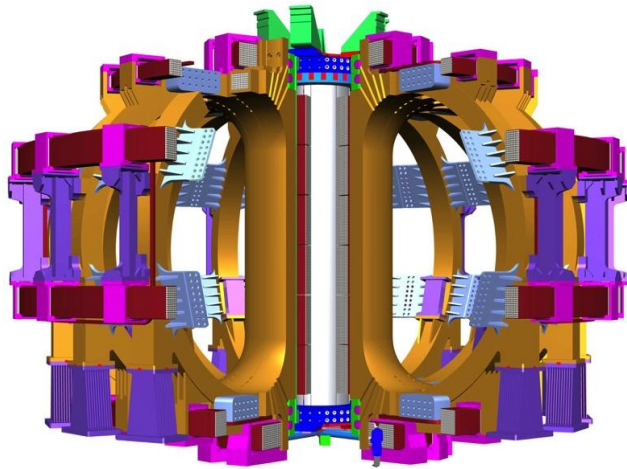
K.Shinohara



EPFL Fast ions and the self-heating process

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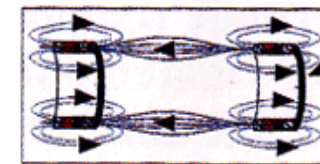
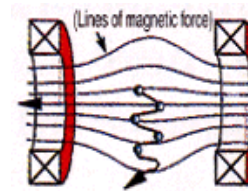
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Ripple effects on fast ions qualitatively understood

Ferritic inserts in ITER to reduce ripple losses (from ~6% to ~0.4%)



Ferritic steel
(Ferromagnetic material)

K.Shinohara

Open questions

Effect of blanket modules (~1% ripple)

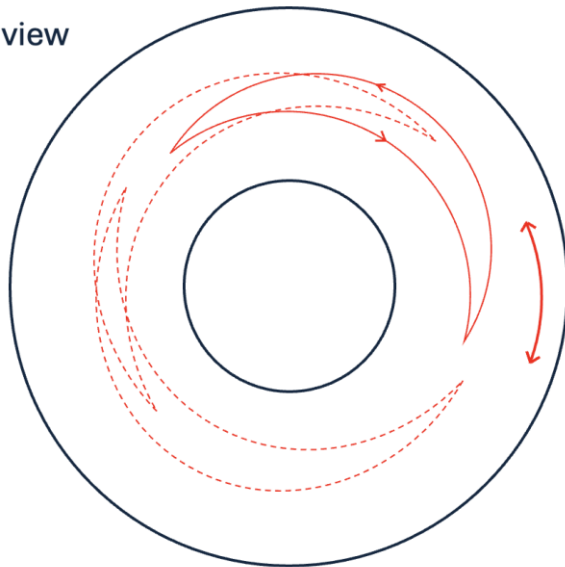
Effect of 3D fields induced by coils introduced to mitigate edge instabilities

FAST ION LOSS MECHANISMS

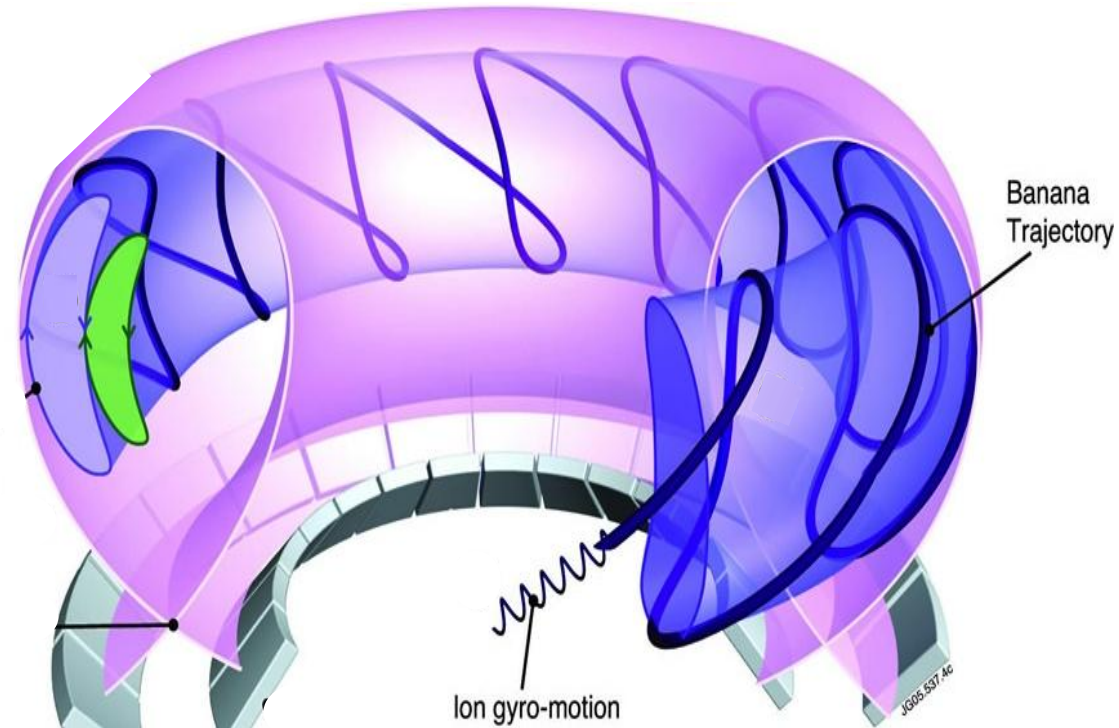
LOW FREQUENCY MHD

Low frequency MHD Resonance condition

Top view



ω
precession



Ion gyro-motion

Banana
Trajectory

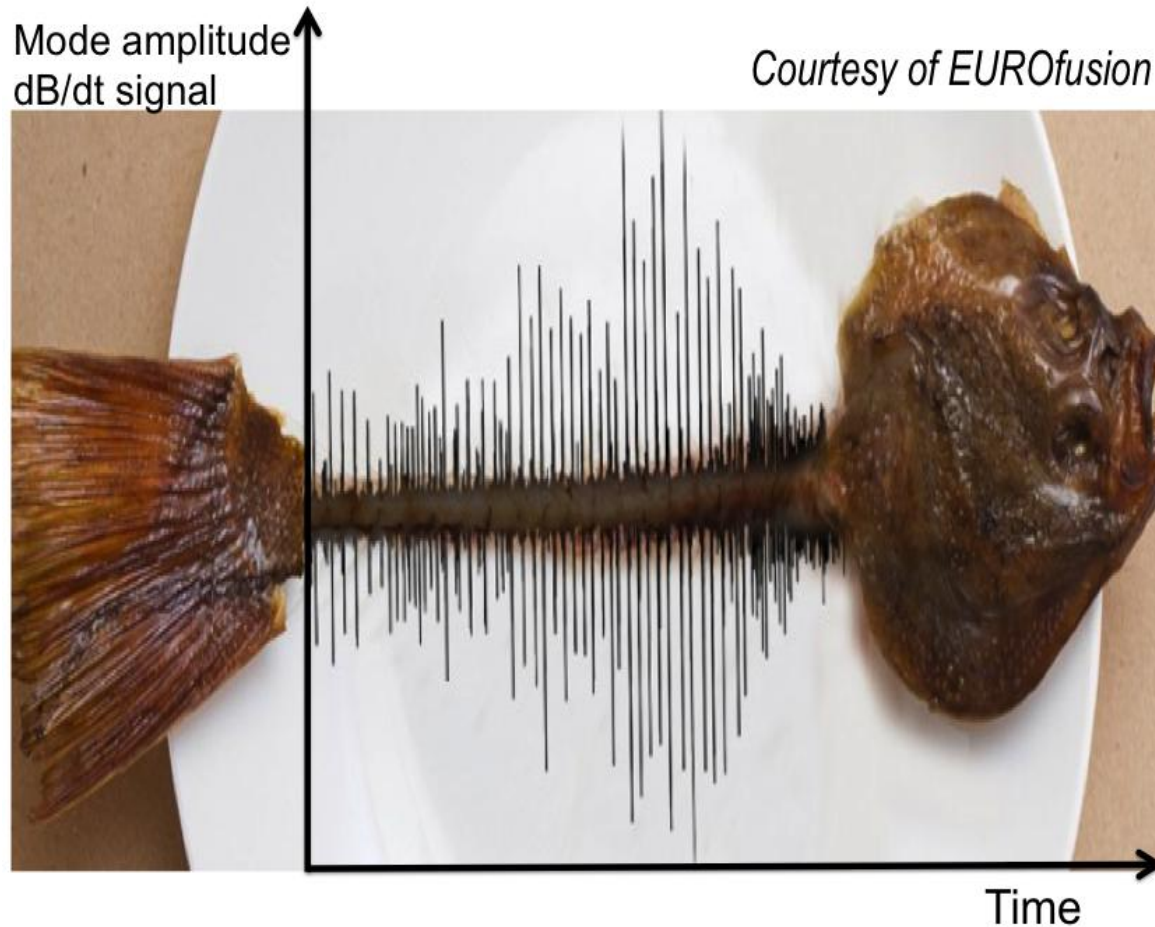
Different kinds of resonances associated with complex particle orbits lead to different possible interactions

Example: bounce motion of particles along banana orbits and toroidal drift precession of banana orbits : $\omega_{\text{precession}} \ll \omega_{\text{bounce}} \ll \Omega_c$

MHD modes: $\omega \ll \Omega_c$

EPFL Low frequency MHD – fishbones

Resonant de-stabilisation of MHD kink mode with $\omega = \omega_{\text{precession, fast ions}}$
Driven by the fast ions, the mode can reach amplitudes that cause ejection of the fast ions themselves - the mode then disappears as its source is gone - a sequence of bursts can occur

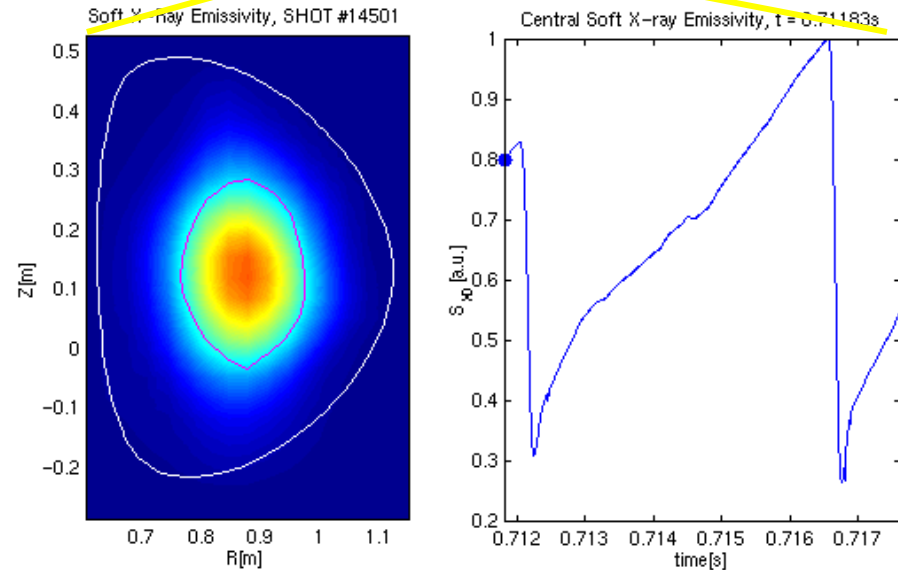
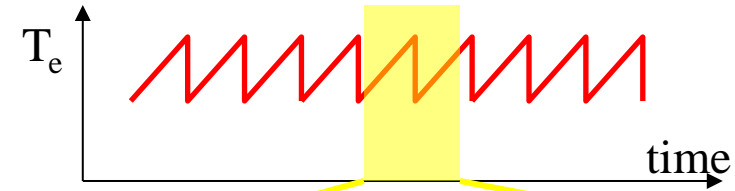


Low frequency MHD Sawtooth instability

Sudden losses of energy and particles in core, with local breaking of magnetic structure (magnetic reconnection)

Kink mode with $\omega < \omega_{\text{precession}}$, fast ions
Fast ions are not the cause of the instability, but are ejected together with core plasma

Ex.: X-ray emissivity ($\sim T_e, n$) evolution in TCV



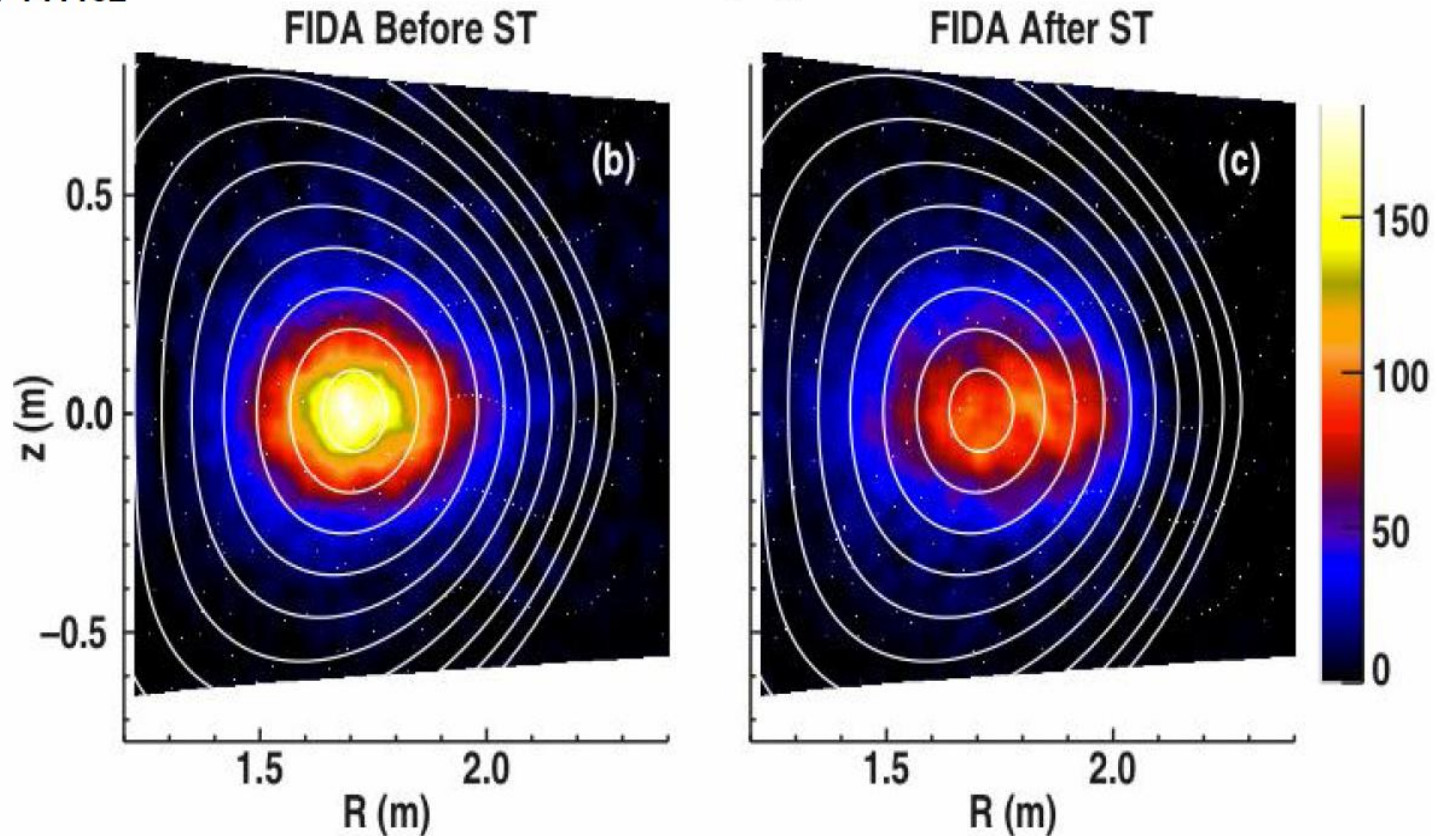
Courtesy of I. Furno

Collapses can trigger secondary instabilities: as fast ions influence the sawtooth period, hence the strength of the collapse, they can be used to control the secondary instabilities

Low frequency MHD Sawtooth instability

Redistribution measured by Fast ion $D\alpha$

141182



FAST ION LOSS MECHANISMS

PLASMA TURBULENCE

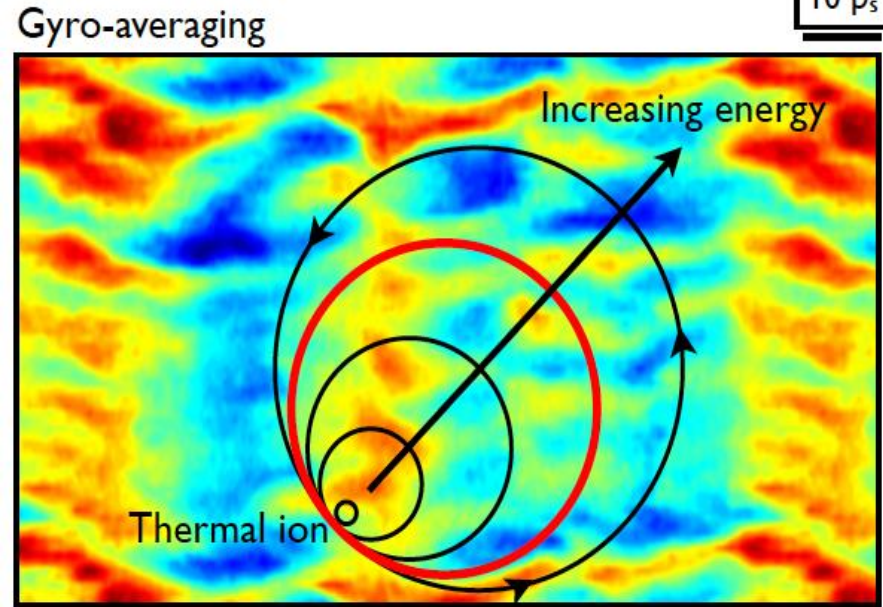
EPFL Fast ion interaction with turbulence

$10 \rho_s$

Turbulence could cause transport of fast ions, as for thermal plasma

Large fast ion orbits are expected to average out effect of turbulence

M. Albergante



Key parameters: $E_{\text{fast ions}}/T_{\text{plasma}}$ & fast ion slowing down time

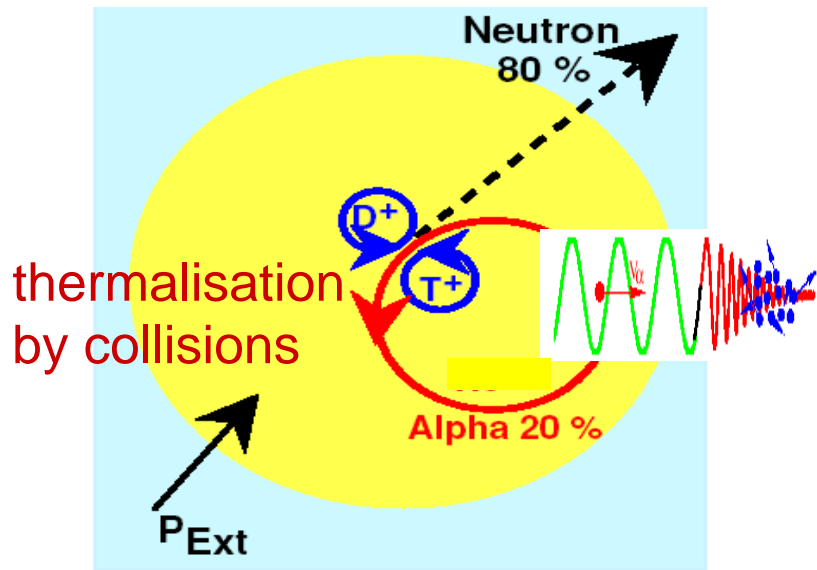
Present devices (small $E_{\text{fast ions}}/T_{\text{plasma}}$):
some anomalous transport of NBI ions

ITER (large $E_{\text{fast ions}}/T_{\text{plasma}}$):
negligible effect for NBI ions and for α 's

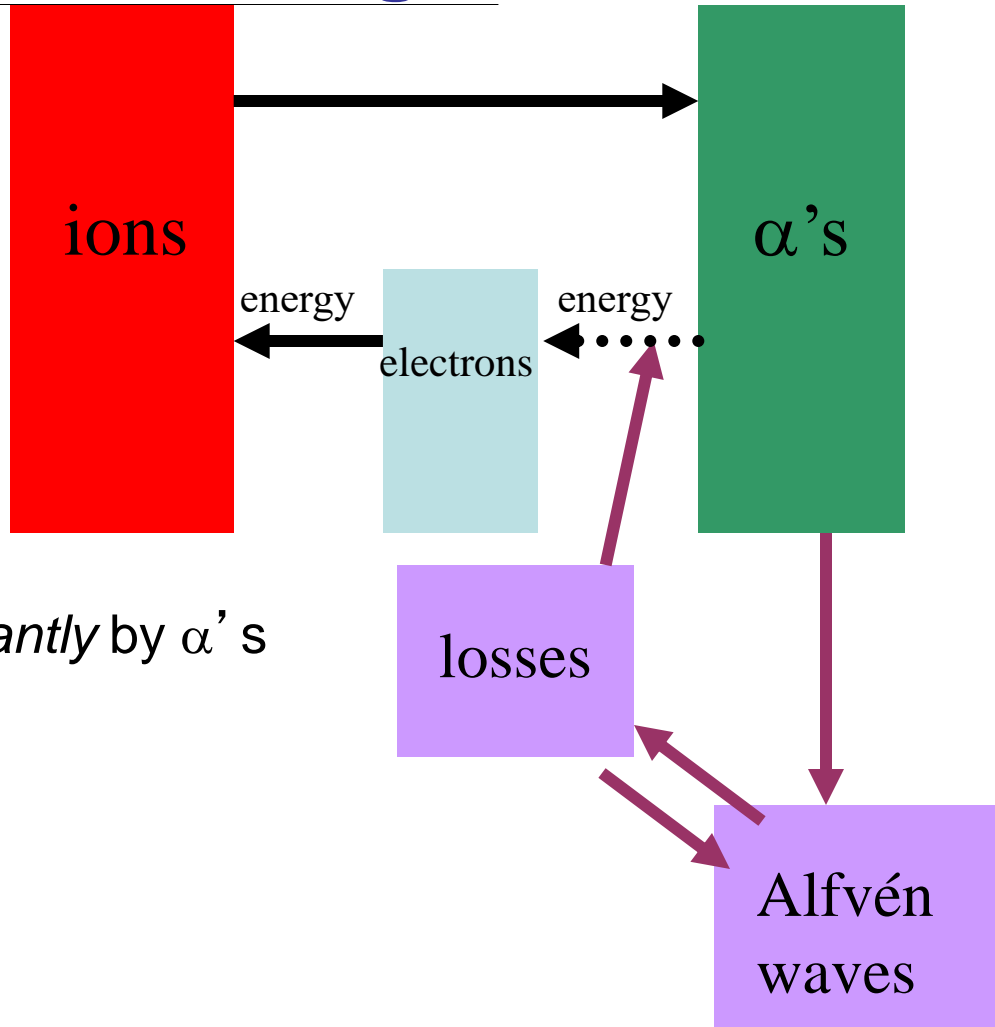
FAST ION LOSS MECHANISMS

INTERACTION WITH ALFVEN WAVES

Resonant interaction with waves and self-heating



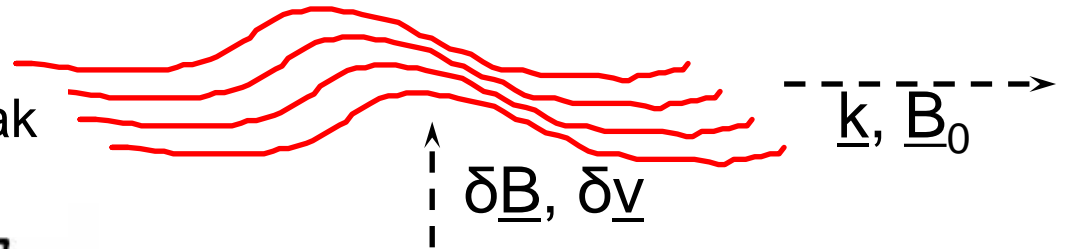
Alfvén waves can be driven *resonantly* by α 's



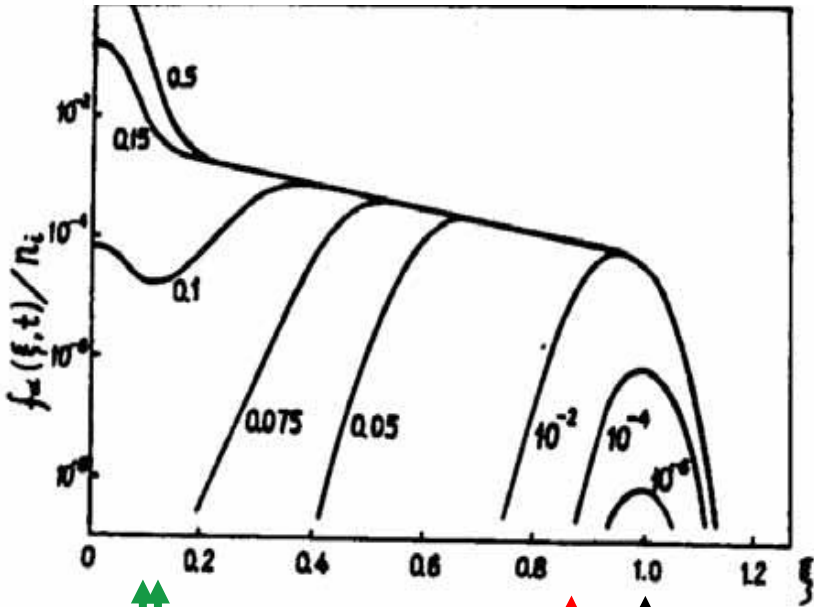
Fast ions and Alfvén waves

B-field and plasma frozen together; field lines are strings with tension and inertia → Alfvén wave propagation

Typical velocities in a tokamak
 $B=4T$; $T=10keV$; $n=10^{20}m^{-3}$

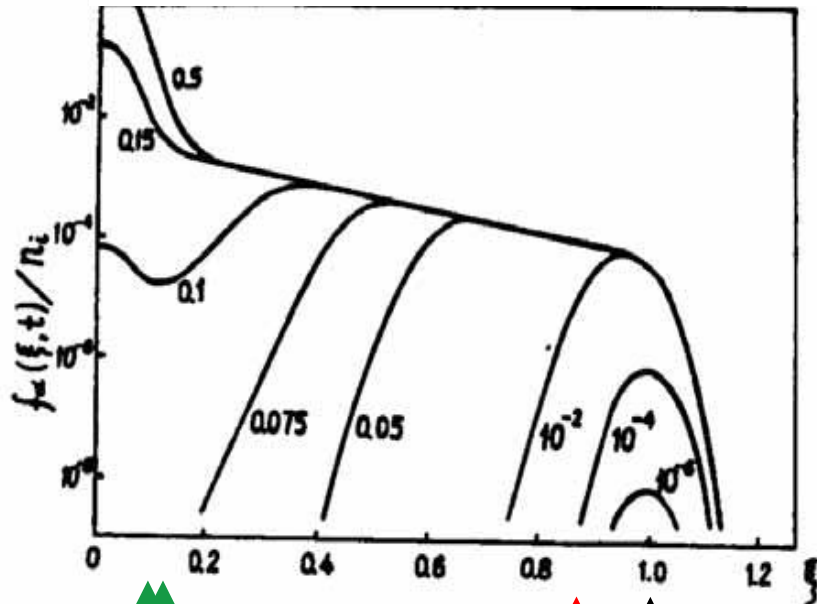


Slowing down α 's (but also fast ions generated by additional heating) can resonate with Alfvén Waves



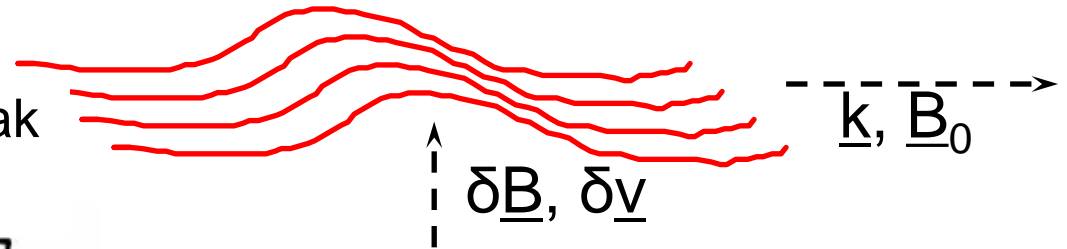
B-field and plasma frozen together; field lines are strings with tension and inertia \rightarrow Alfvén wave propagation

Typical velocities in a tokamak
 $B=4\text{T}$; $T=10\text{keV}$; $n=10^{20}\text{m}^{-3}$



V_{thermal} \uparrow D, T

$V_{\text{Alfvén}}$ \uparrow $V_{\text{birth}}^{\alpha}$



Slowing down α 's (but also fast ions generated by additional heating) can resonate with Alfvén Waves

Alfvén Waves are driven unstable if the 'free' energy ∇p_{α} is sufficient and α drive $>$ plasma damping

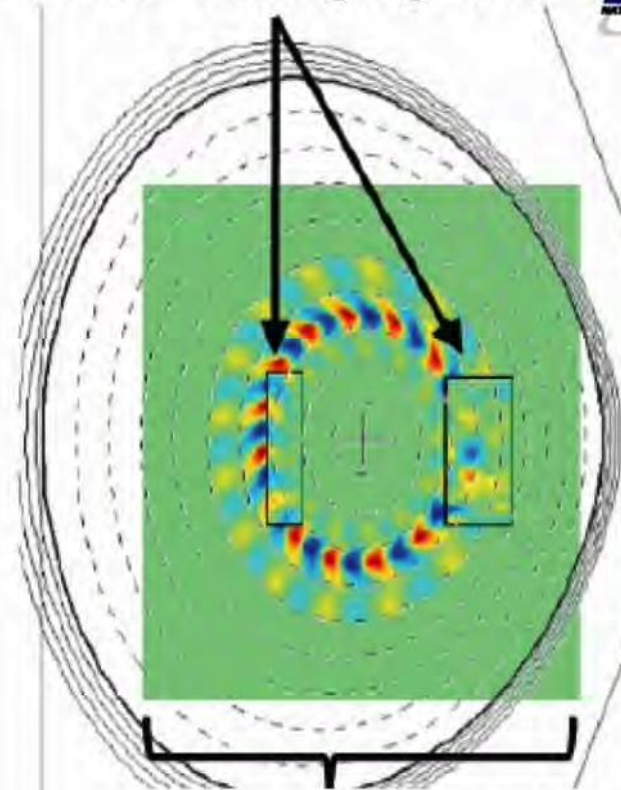
Alfvén waves in tokamaks

Alfvén Eigenmodes

Alfvén waves' dispersion in tokamaks allows for weakly damped global Eigenmodes

MHD theory has successfully predicted the existence and the main features of these modes, then verified experimentally

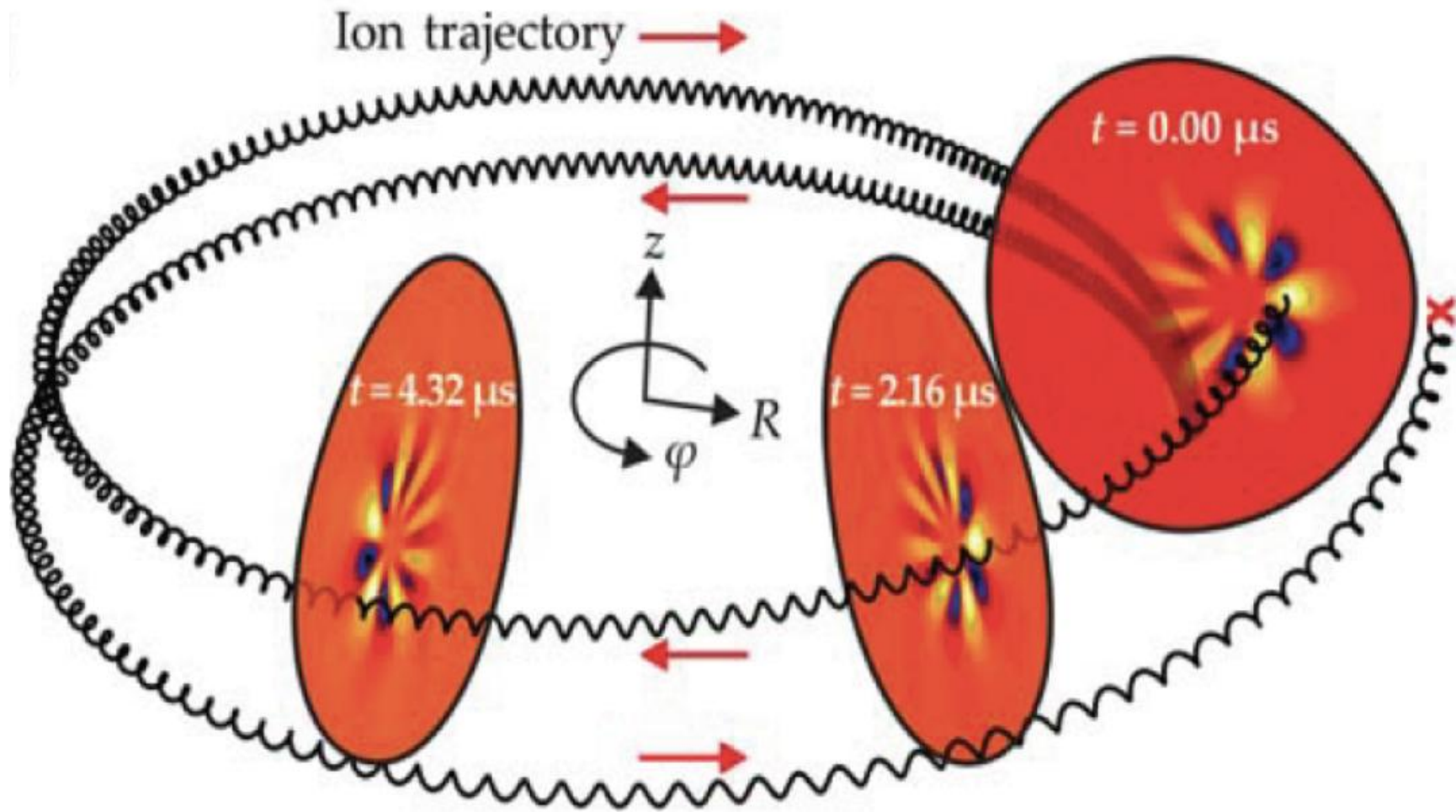
Electron cyclotron emission imaging data



Calculated structure of
Alfvén Eigenmode

Courtesy of B.J. Tobias

AE resonance condition visualisation



EPFL Fast ions and the self-heating process

Effect of Alfvén Eigenmodes

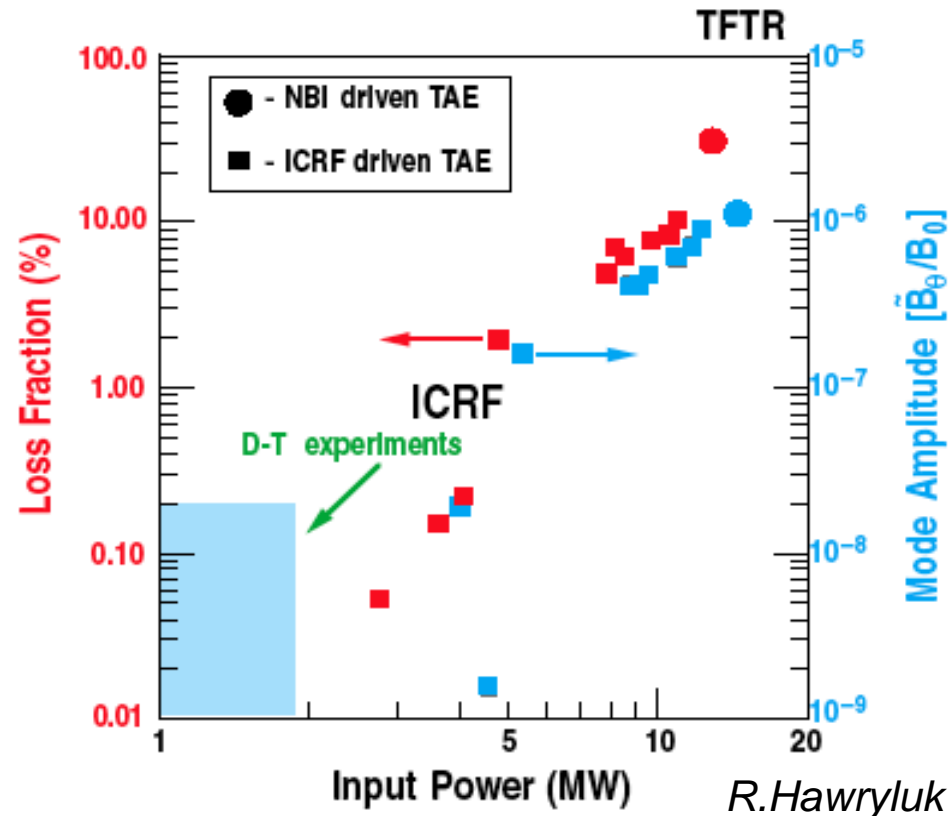
Losses and redistribution seen in many experiments

ITER / DEMO can withstand only a few % of α losses

Questions

Linear stability, balance between damping and drive

Nonlinear evolution and redistribution / losses

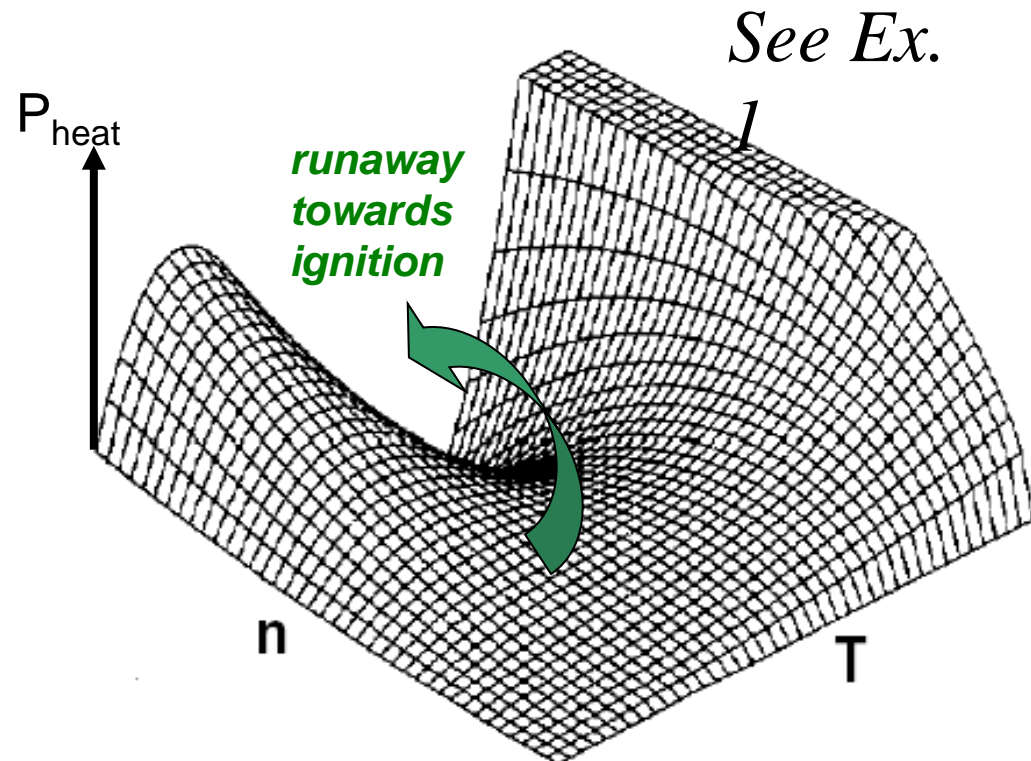
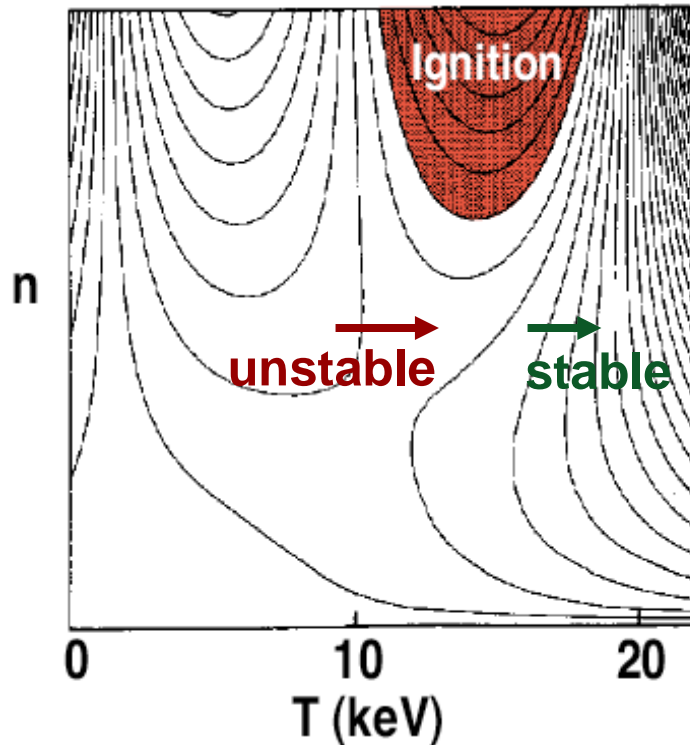


BURN STABILITY AND CONTROL

Burn thermal runaway

At high Q in principle thermal runaway can occur

$$\text{Ex. of steady-state } P_{\text{in}}/V = \underbrace{3nT/\tau_E}_{P_{\text{loss}}} - \underbrace{1/4n^2\langle\sigma v\rangle E_\alpha}_{P_\alpha}$$



TRITIUM IN FUSION

Half life time: 12.3 years

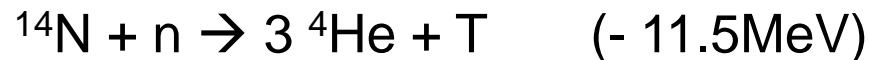
Decay mode

β decay: $T \rightarrow {}^3\text{He} + e^- + \bar{\nu}_e + 18.6 \text{ keV}$

Activity = 3.6×10^{14} Bq/g

Tritium – natural inventory

~3.6 kg produced mostly by impact of cosmogenic neutrons on Nitrogen



~30 kg left from atmospheric testing of nuclear weapons (1945-1963)

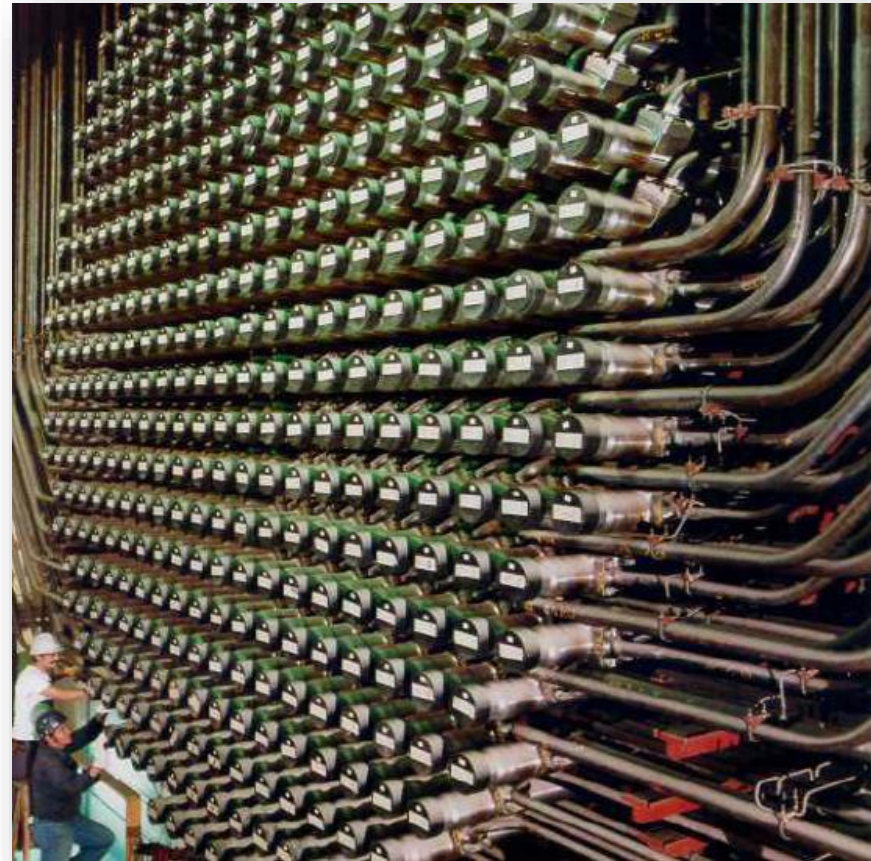
Tritium – anthropogenic inventory

T is generated mainly in heavy water fission reactors of CANDU type (CANadian Deuterium natural Uranium reactor)

2003: 20 operating CANDU reactors
(40 years license)

Total production
 $\sim 0.3\text{kg}/(\text{GWe} \cdot \text{year})$

Today 19kg at hands in
Darlington recover facility



CANDU Pressurized Heavy Water Reactor

Tritium production is today driven by fission (CANDU) waste management rather than by a market



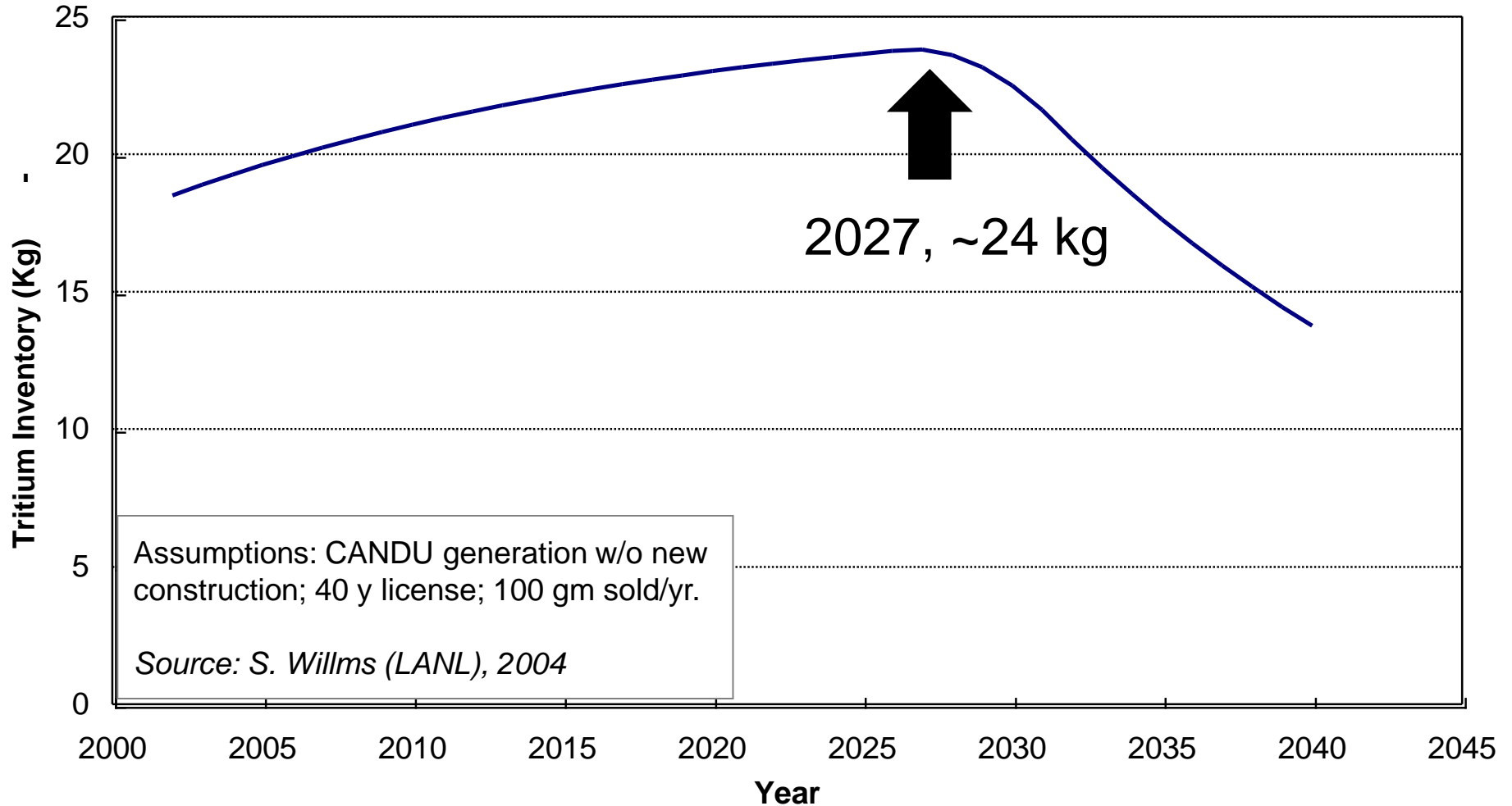
Darlington (Canada)
Tritium Removal Facility



~1.5kg tritium recovered/year

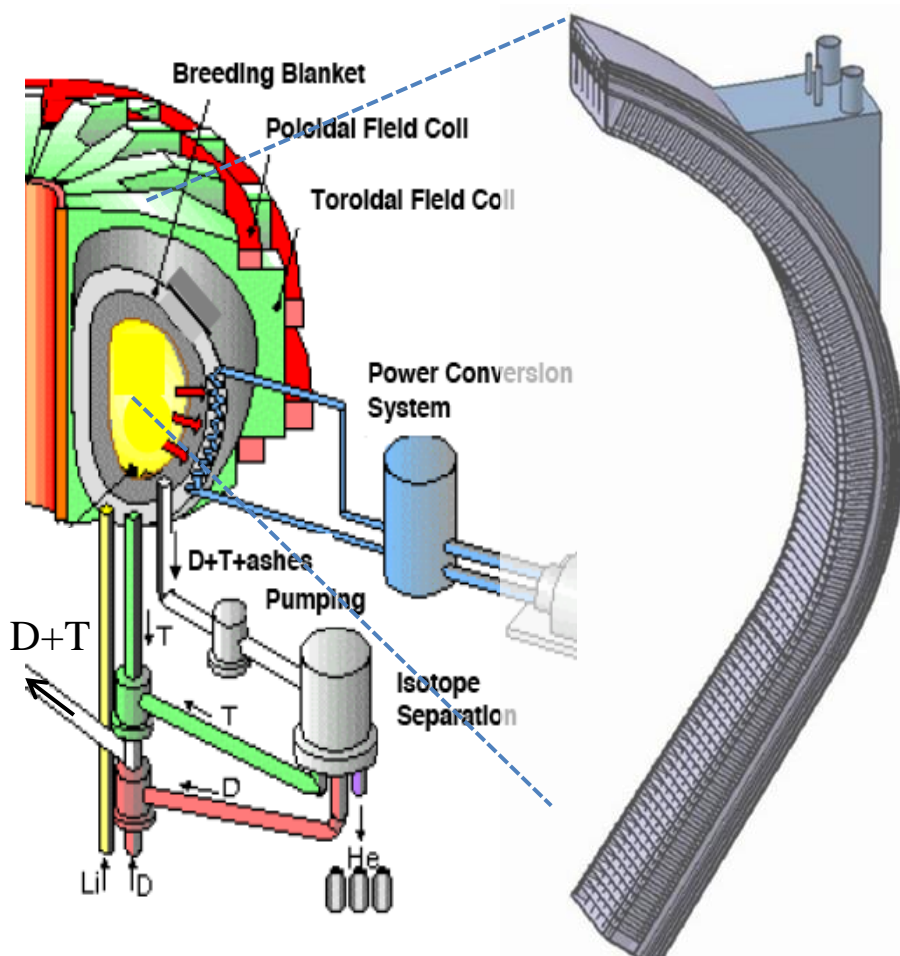
EPFL T availability is foreseen to decrease after 2027

Projection of tritium inventory available from CANDU (Canada)



But other countries (e.g. India) also build heavy water reactors

Tritium breeding blanket



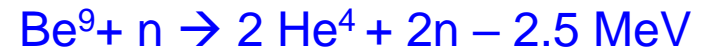
A fusion reactor must be T self sufficient



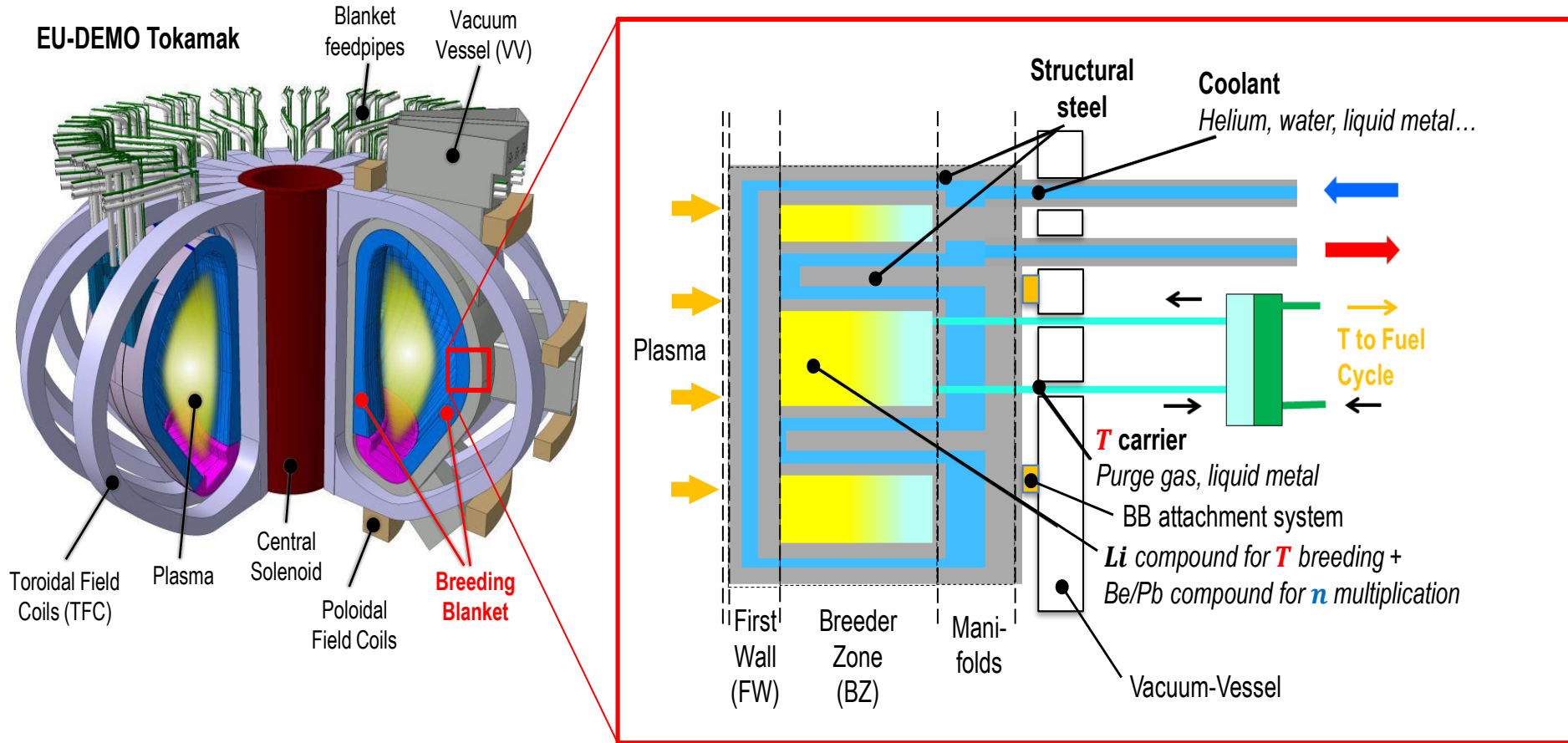
Tritium breeding ratio

TBR = tritium bred / tritium burnt

TBR > 1 to compensate losses, which implies the use of n-multiplier



EPFL Breeding blanket – system and functions

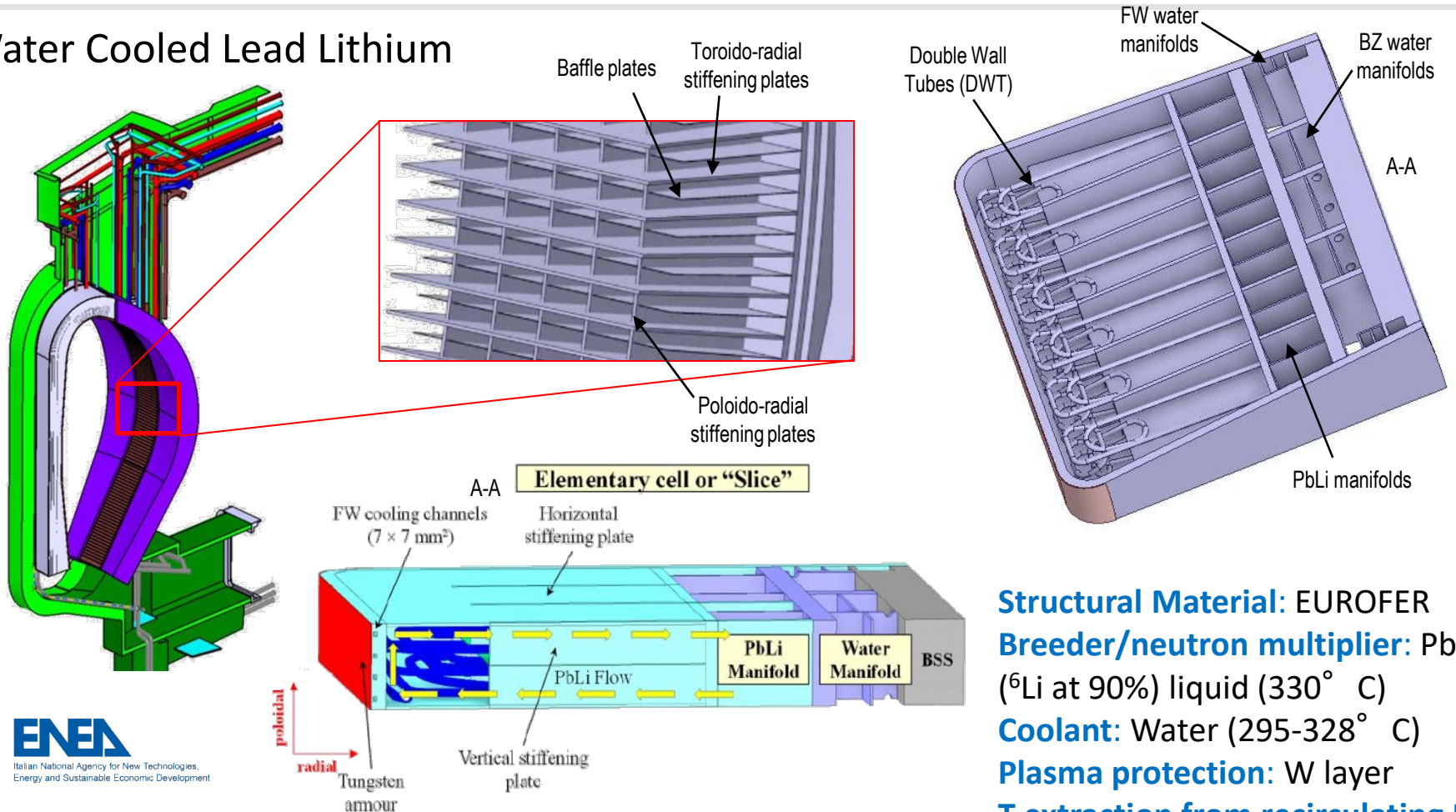


F.A. Hernández et al. | BB Functions, Concepts and Associated Issues | Fusion/Fission Workshop, VC | 06.12.2023 | Page 4

Several concepts are studied ...

EPFL Ex. of BB concept for DEMO - WCLL

Water Cooled Lead Lithium

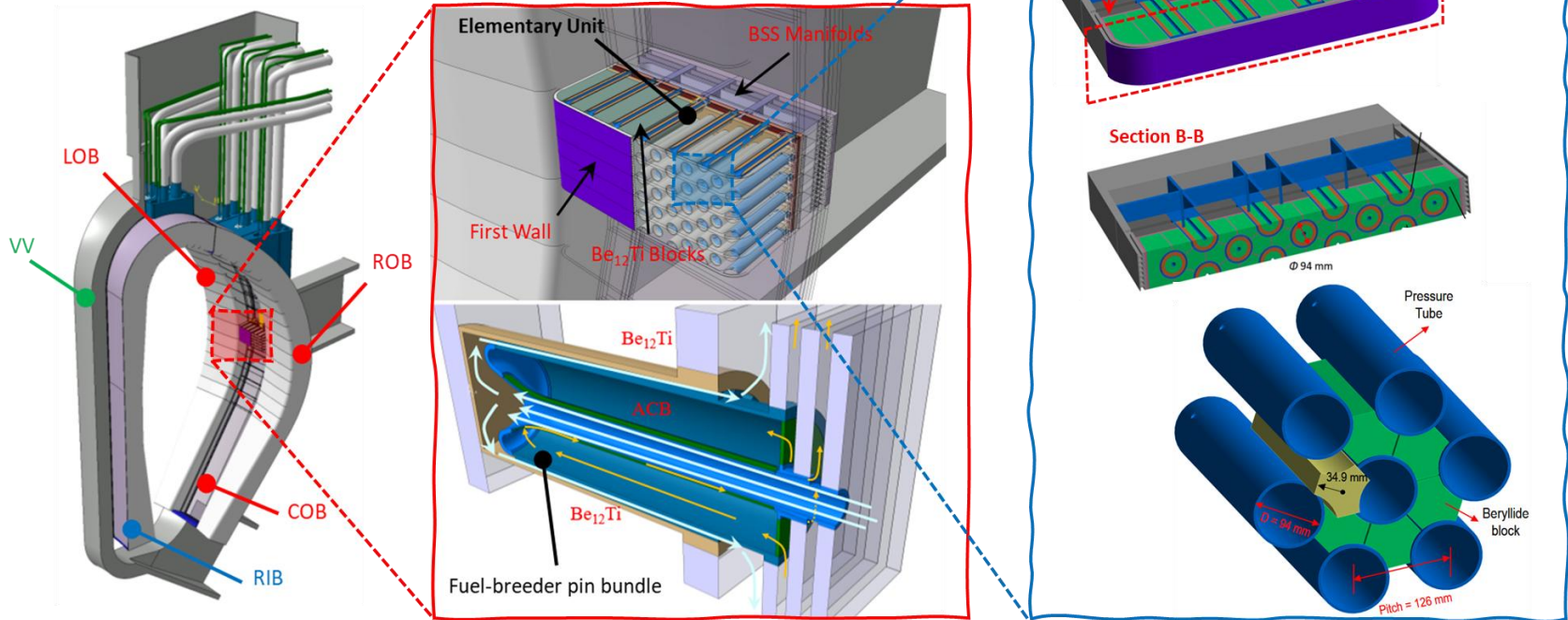


ENEA
Italian National Agency for New Technologies,
Energy and Sustainable Economic Development

P. Arena et al, (2023) Design and Integration of the EU-DEMO Water-Cooled Lead Lithium Breeding Blanket, *Energies*, 16, 2069

EPFL Ex. of BB concept for DEMO - HCPB

Helium Cooled Pebble Bed



Structural Material: EUROFER

Breeder: $\text{Li}_4\text{SiO}_4 + \text{Li}_2\text{TiO}_3$ in form of a pebble bed, ^6Li at 60%

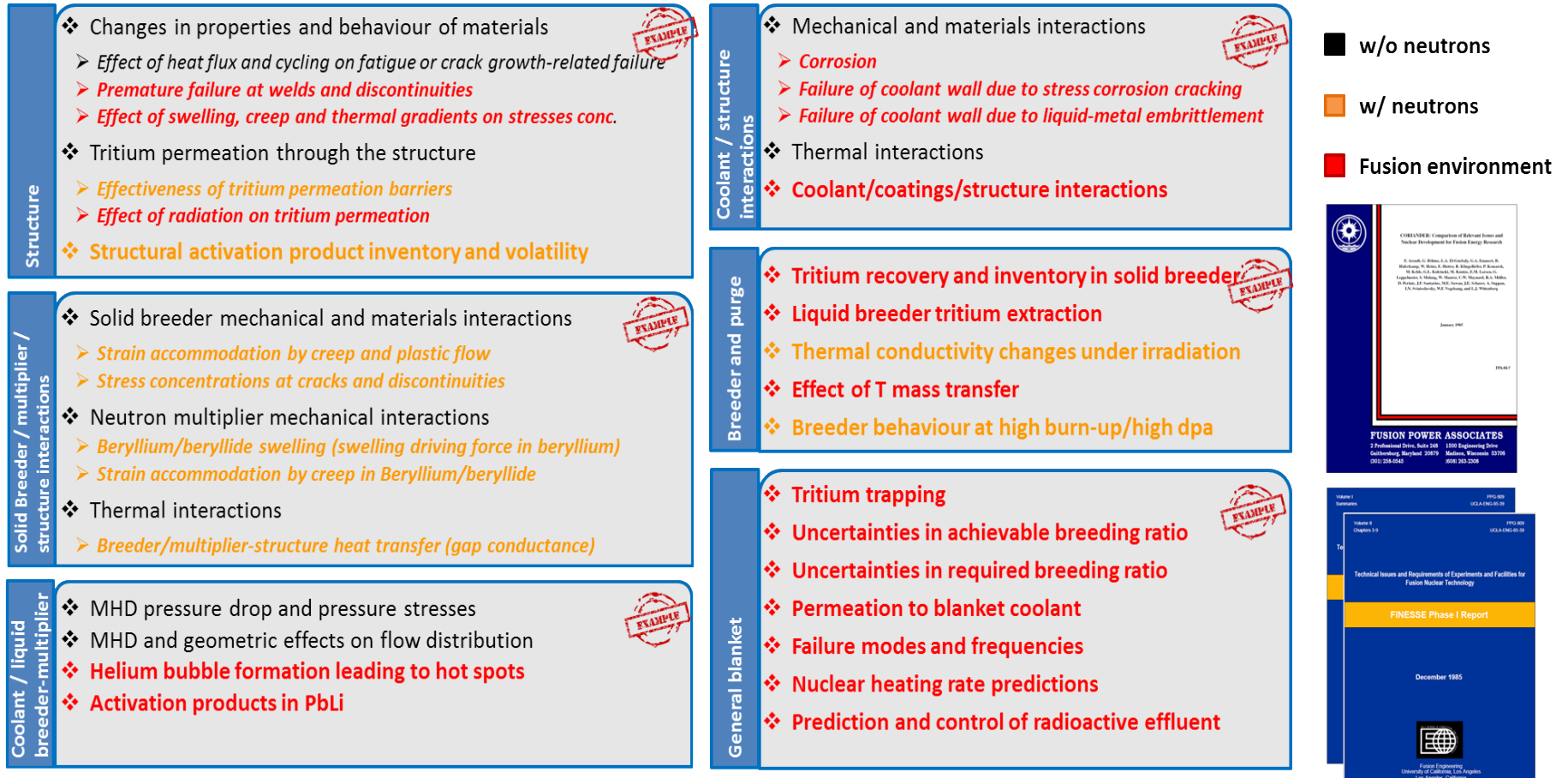
Neutron multiplier: Beryllide (TiBe_{12}) rods

Coolant: He 300-520° C @ 8 MPa

Plasma protection: W layer

T extraction with purge Helium

EPFL Challenges related to the BB



Efforts should be planned to test mock-ups and prototypes under relevant conditions to improve the understanding and knowledge of phenomena/processes. Separate effect tests can be foreseen initially. However, integral testing are needed for a full BB qualification.

BB in DEMO: improved variants

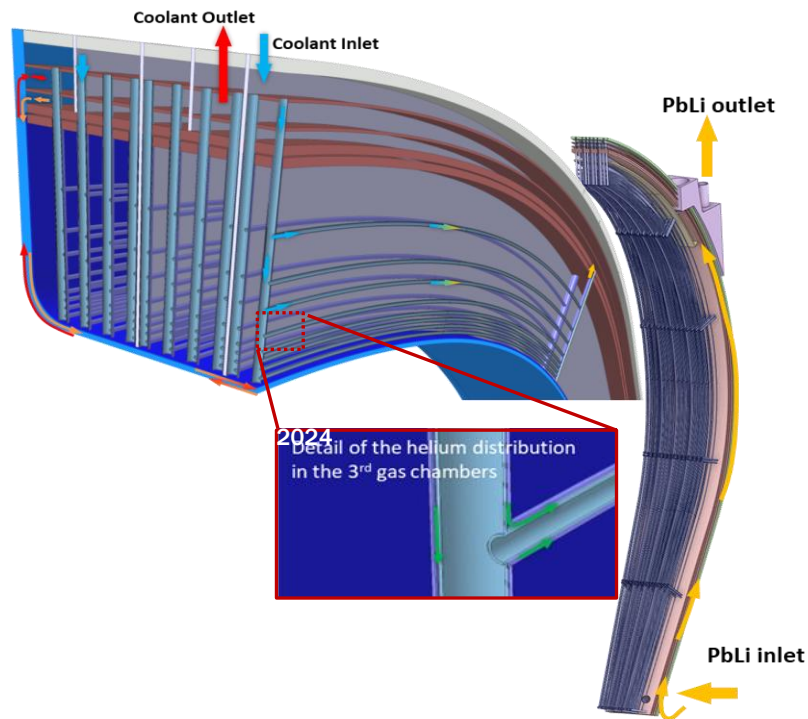
Water Cooled Lead Lithium Concept – Double Bundle

Breeder zone and first wall in series ($T_{in,FW} \approx 315^\circ\text{C}$)

Less tubes, less welds, less surface, less water, more PbLi

→ more reliable, lower T permeation, better TBR

Avoids PbLi - water interaction



Water-cooled Lead and Ceramic Breeder

Neutron multiplier: molten Pb or solid Pb compound

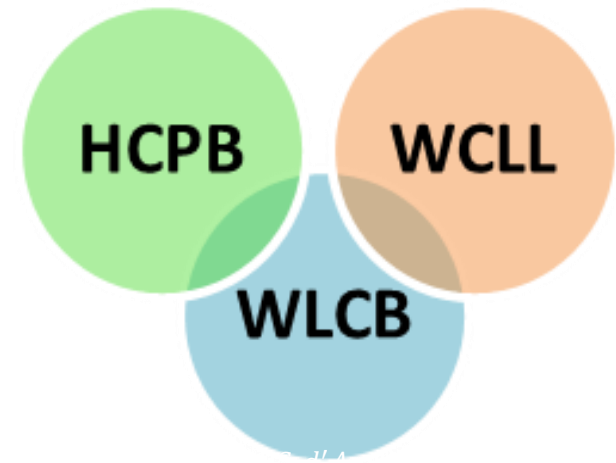
Coolant: Water (285-325°C @15.5 MPa)

T extraction with purge Helium

Breeder zone and first wall in series ($T_{in,FW} \approx 315^\circ\text{C}$)

Available TER and water PWR technology

Good TBR and neutron shielding, low T permeation



Tritium Burn-up Fraction

Tritium Burn-up Fraction f_B = probability that a tritium atom injected into the plasma will undergo a fusion reaction before it escapes

$$f_B = \text{fusion reaction rate} / \text{tritium fueling rate} = n_T^2 \langle \sigma v \rangle_{DT} / S_T$$

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$$f_B = \text{fusion reaction rate} / \text{tritium fueling rate} = n_T^2 \langle \sigma v \rangle_{DT} / S_T$$

T particle balance ($n_T = n_D$):
$$dn_T/dt = S_T - n_T^2 \langle \sigma v \rangle_{DT} - n_T/\tau_T$$

τ_T = particle confinement time - with edge recycling τ_T becomes $\tau_T/(1-R)$,
but generally the recycling coefficient $R \ll 1$

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τ_T = particle confinement time - with edge recycling τ_T becomes $\tau_T/(1-R)$,
but generally the recycling coefficient $R \ll 1$

Steady-state:
$$S_T = n_T^2 \langle \sigma v \rangle_{DT} + n_T/\tau_T$$

and

$$f_B = n_T^2 \langle \sigma v \rangle_{DT} / S_T = n_T^2 \langle \sigma v \rangle_{DT} / (n_T^2 \langle \sigma v \rangle_{DT} + n_T/\tau_T) = 1 / (1 + 1/(\tau_T n_T \langle \sigma v \rangle_{DT}))$$

f_B can be increased by increasing the confinement time

EPFL Tritium inventory in fusion reactor

dM/dt = mass consumption rate to produce the necessary fusion power

$$dM/dt \sim 56 \times \text{fusion power kg/y/GW}_{\text{thermal}}$$

$$\text{Ex. } 1\text{GW}_e \rightarrow \sim 3\text{GW}_{\text{thermal}} \rightarrow \sim 160\text{kg/y}$$

T reprocessing system takes a mean time t_p to clean up and recycle tritium

The reprocessed tritium is injected into the plasma with efficiency η_f

γ_s = radioactive decay rate ($= \ln 2 / 12.3 \text{ y}^{-1}$); γ_r = loss rate in reprocessing T

M_0 = time-independent, re-circulating tritium inventory

EPFL Tritium inventory in fusion reactor

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Over a time interval t_p

$$\underbrace{M_0 \times \underbrace{h_f}_{\text{injected}} \times f_B}_{\text{burnt}} \times \underbrace{TBR}_{\text{bred}} = \int_0^{t_p} \left[\frac{dM}{dt} + (g_r + g_s)M_0 \right] dt = \left[\frac{dM}{dt} + (g_r + g_s)M_0 \right] \times t_p$$

$$dM/dt = \eta_f f_B M_0 / t_p \times TBR - (\gamma_r + \gamma_s) M_0$$

EPFL Tritium inventory in fusion reactor

$$dM/dt = \eta_f f_B M_0 / t_p \times TBR - (\gamma_r + \gamma_s) M_0$$

If breeding rate is large compared to loss rate, and $TBR \sim 1$

$$M_0 \sim t_p dM/dt / \eta_f f_B$$

Ex.: 1 GW_{thermal} reactor, with $f_B = 5\%$, $\eta_f = 50\%$, $t_p = 1\text{day} \rightarrow M_0 \sim 6\text{kg}$
ITER $\rightarrow M_0 \sim 3\text{kg}$ ($\sim \$100\text{M}$)
DEMO $\rightarrow M_0 \sim 7\text{kg}$

Naturally, the un-burnt T is also recycled, with the same processing time t_p

Problems with large T inventory

Safety; Power required to heat fueled T to plasma temperature; Required TBR

T inventory can be minimized by minimizing the reprocessing time and by increasing the injection efficiency and burn-up fraction

EPFL Tritium build-up in fusion reactors

In reality, the total tritium mass in a reactor should increase with time

$M = M_0 + m$, with m the mass produced in reactor blanket

$$\frac{dm}{dt} = -g_s m - [g_s + g_r] M_0 - \frac{dM}{dt} + \frac{dM}{dt} TBR$$

As $dM/dt \sim M_0 \eta_f f_B / t_p$

$$\frac{dm}{dt} = -g_s m + \underbrace{\frac{\eta_f f_B}{t_p} (TBR - 1) - g_s - g_r}_{A} M_0 = -g_s m + AM_0$$

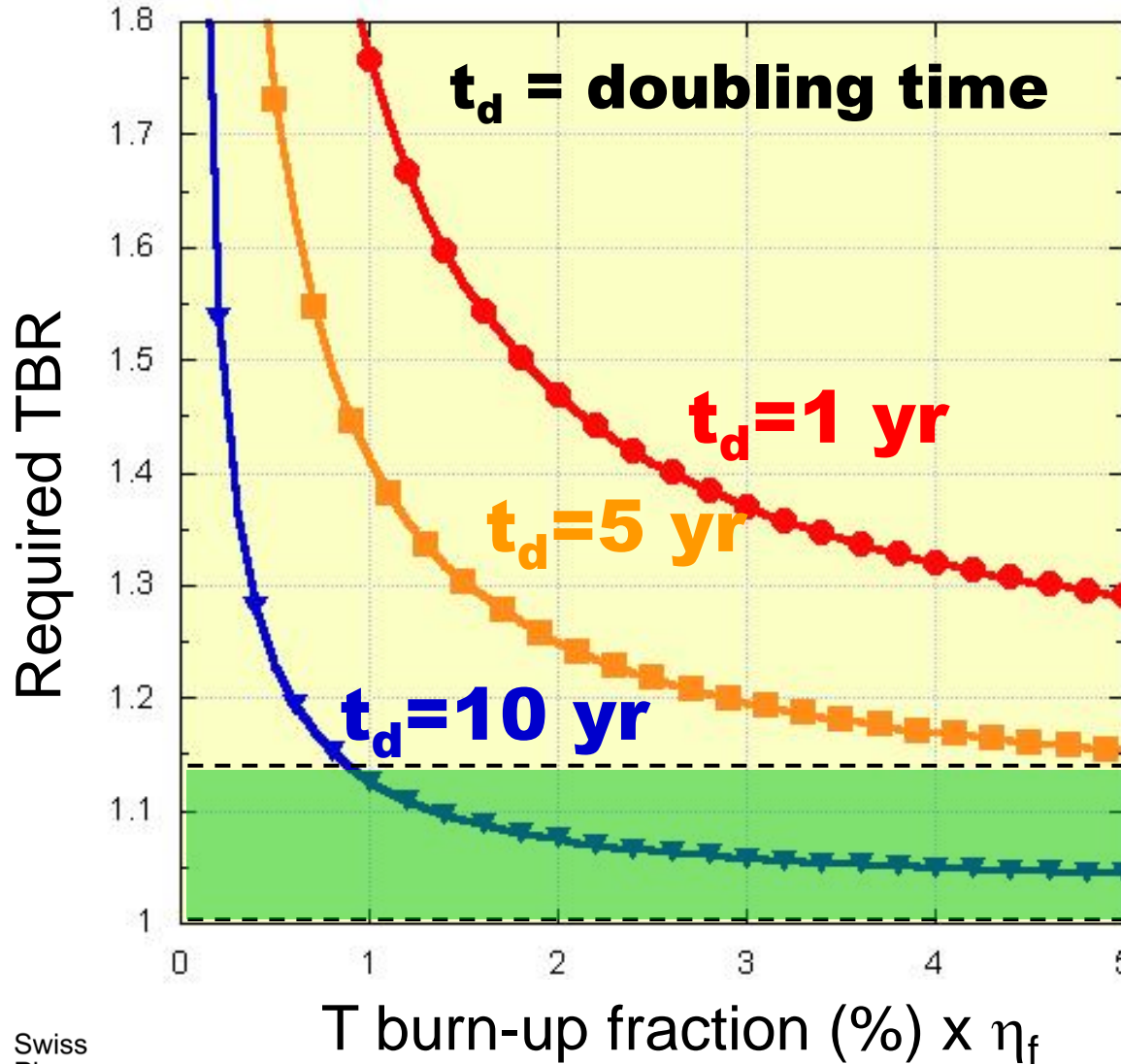
Which gives

$$m(t) = \frac{AM_0}{g_s} (1 - e^{-g_s t})$$

This regulates the 'tritium economy' and tells us how fast we can deploy reactors

Impact of T cycle parameters on fusion

Required TBR for T self-sufficiency depends on burn-up fraction, fuelling efficiency and fusion reactor doubling time



Fusion power ~ 1.5 GW
 Waste removal efficiency ~ 0.9
 (from Prof. Abdou, UCLA)

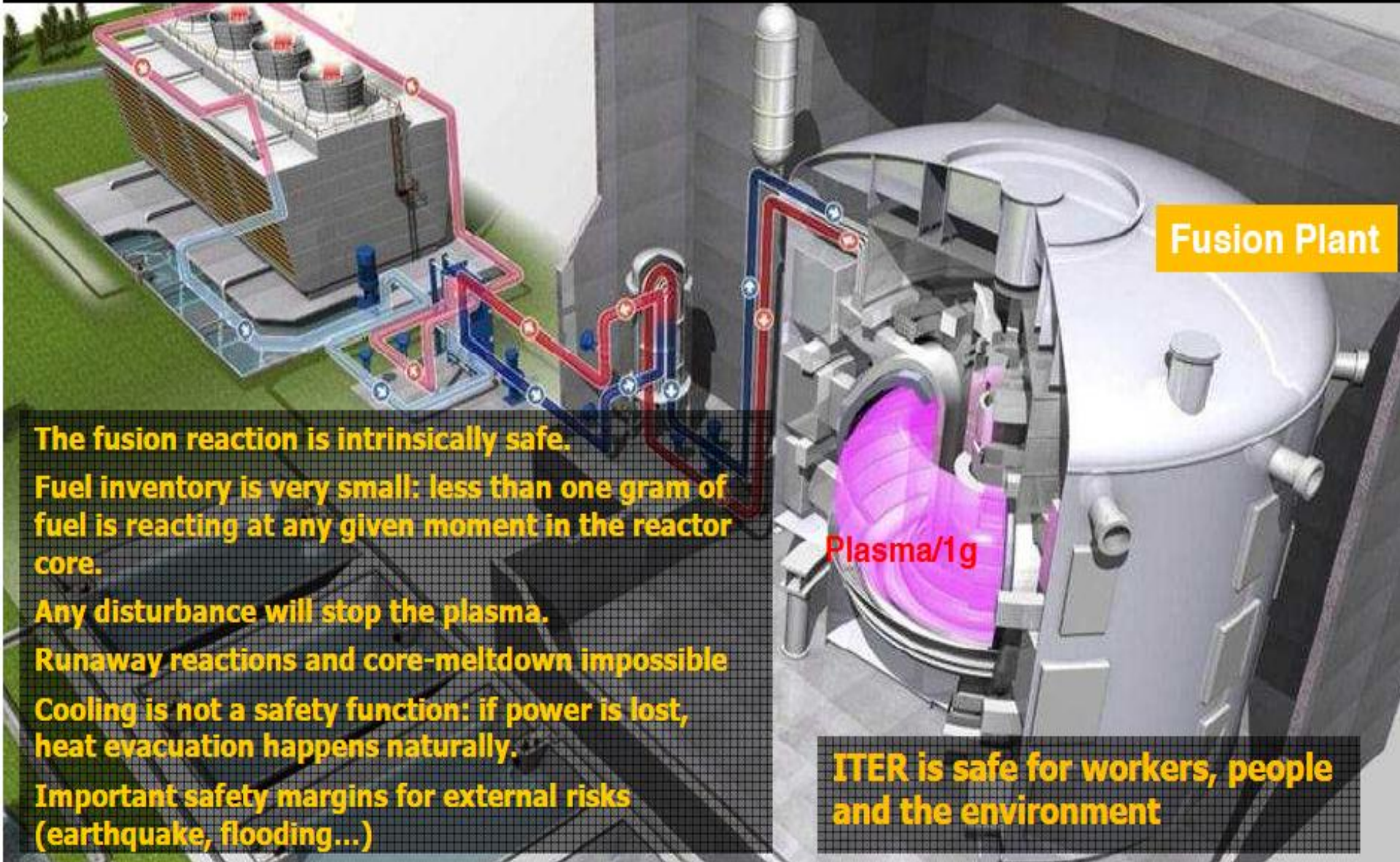
Max TBR ≤ 1.15

"Window" for Tritium self sufficiency

Significant R&D needed to stay in window

ITER SAFETY AND LICENSING

A Fukushima-like accident is impossible in ITER



The fusion reaction is intrinsically safe.

Fuel inventory is very small: less than one gram of fuel is reacting at any given moment in the reactor core.

Any disturbance will stop the plasma.

Runaway reactions and core-meltdown impossible

Cooling is not a safety function: if power is lost, heat evacuation happens naturally.

Important safety margins for external risks (earthquake, flooding...)

ITER is safe for workers, people and the environment

Temperature in case of loss of cooling

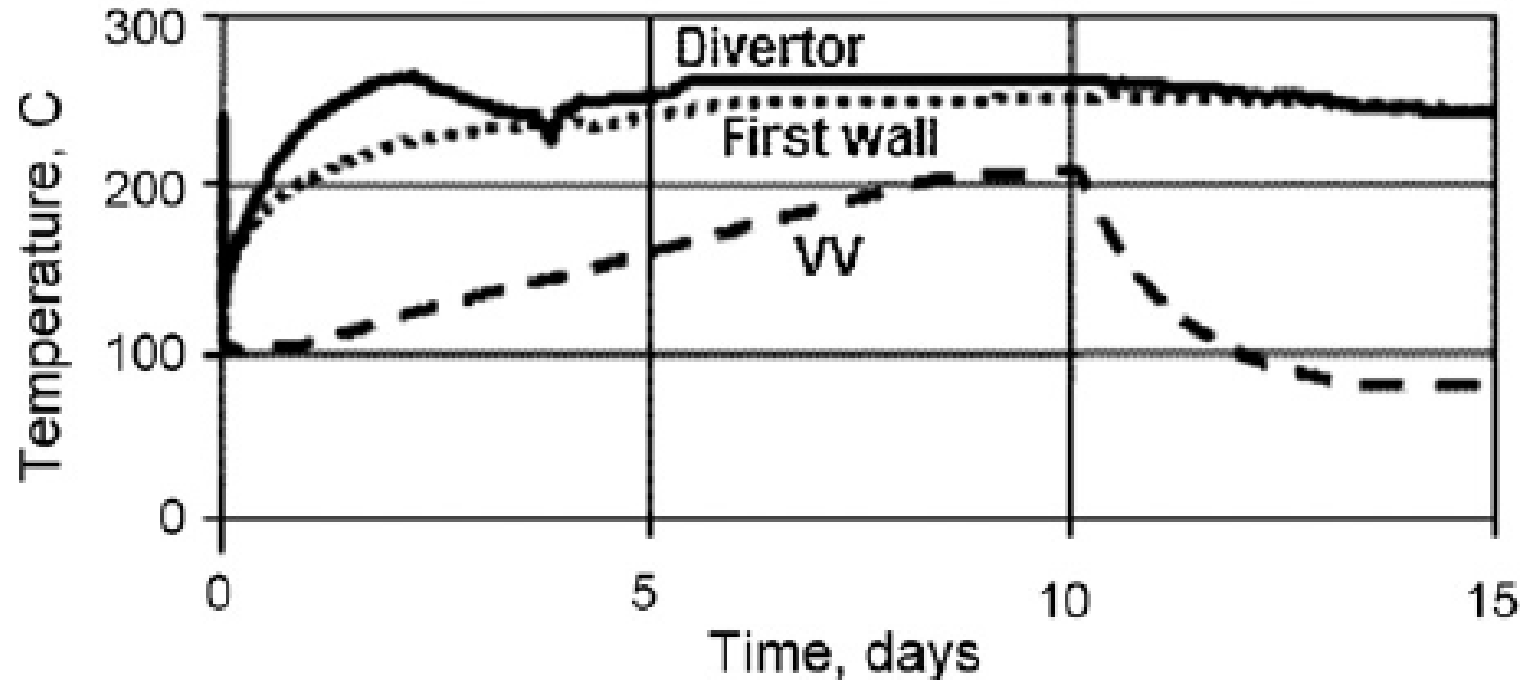


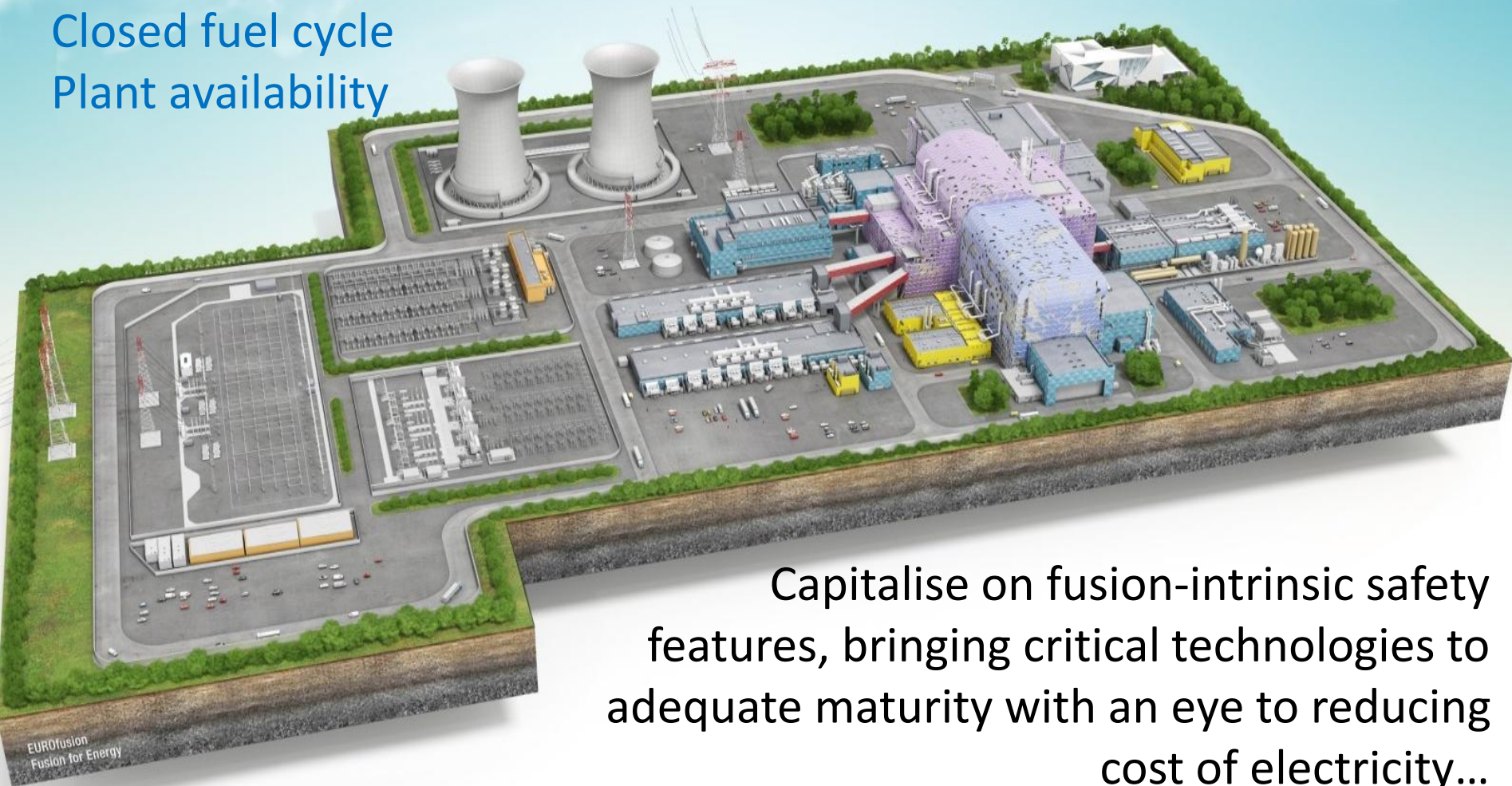
Fig. 1. Temperature evolution of divertor, first wall and vacuum vessel in event of loss of all water coolant flow, with cryostat vented to air after 10 days.

ITER radioactivity and waste

- **ITER will not generate long-lived/high activity waste.**
- **During normal operation, ITER's radiological impact on the most exposed populations will be one thousand times less than natural background radiation.**
- **"Worst-case scenarios", such as fire in the Tritium Plant, would have a lesser impact on neighbouring populations than natural background radiation.**
- **ITER shall observe French safety and security regulations**
 - **A stress test will be conducted by Nuclear Safety Authority.**

TOWARD A FUSION POWER PLANT

Hundreds of MW of net electrical power
Closed fuel cycle
Plant availability



Capitalise on fusion-intrinsic safety features, bringing critical technologies to adequate maturity with an eye to reducing cost of electricity...

Definition of the DEMO step – high level goals

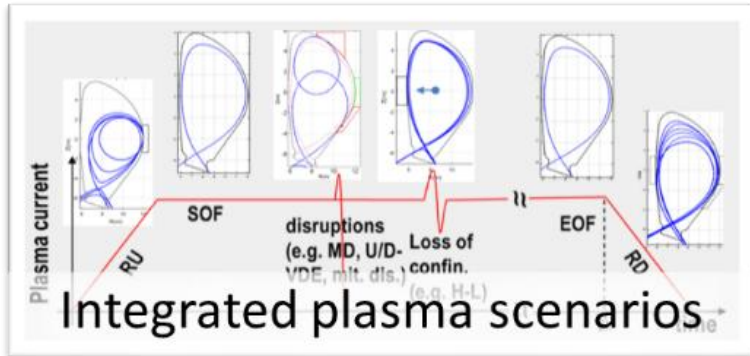
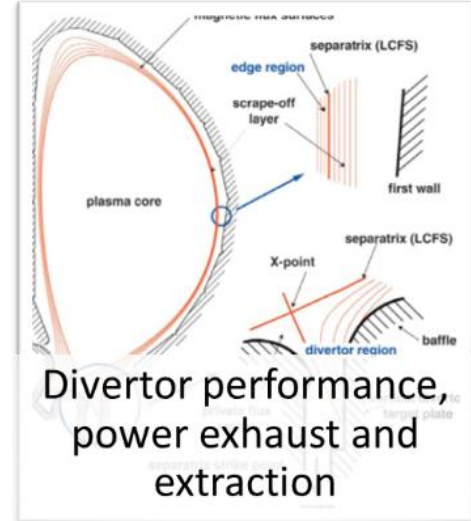
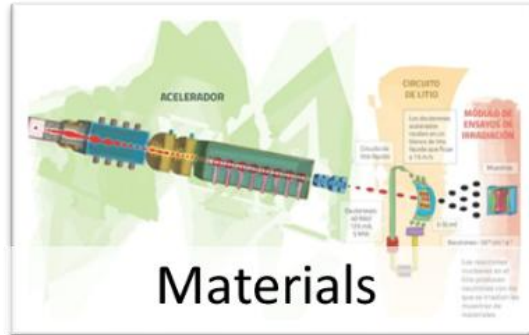
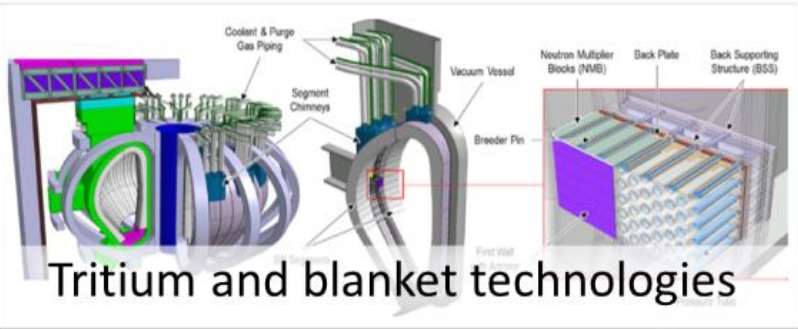
Demonstrate performance and integration of key technologies with tolerable failure rates to achieve adequate levels of availability

- Self-sufficient fuel cycle
- Robust plasma operation scenario & power-exhaust
- Intrinsic safety and tolerable impact of waste
- Phase 1: blanket radiation exposure 20dpa ($\sim 7y$)
- Phase 2: net electricity output $\sim 300MW$ ($t_{pulse} \sim \text{hours}$)
- Tokamak configuration

A.Fasoli | Frascati | November 2024



EPFL Remaining open gaps towards DEMO



EPFL The European Roadmap to fusion

